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Key Points:

- Electron inflow velocities are determined for reconnection at the magnetopause and in the magnetosheath
- For four events inflow velocities of 200–400 km/s imply normalized reconnection rates of 0.05–0.25
- Reconnection rates using electron inflow velocities (0.12) and the reconnection electric field (0.04) are compared for one event

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Electron Inflow Velocities and Reconnection Rates at Earth's Magnetopause and Magnetosheath

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Abstract Electron inflow and outflow velocities during magnetic reconnection at and near the dayside magnetopause are measured using satellites from NASA's Magnetospheric Multiscale (MMS) mission. A case study is examined in detail, and three other events with similar behavior are shown, with one of them being a recently published electron-only reconnection event in the magnetosheath. The measured inflow speeds of 200–400 km/s imply dimensionless reconnection rates of 0.05–0.25 when normalized to the relevant electron Alfvén speed, which are within the range of expectations. The outflow speeds are about 1.5–3 times the inflow speeds, which is consistent with theoretical predictions of the aspect ratio of the inner electron diffusion region. A reconnection rate of $0.04 \pm 25\%$ was obtained for the case study event using the reconnection electric field as compared to the $0.12 \pm 20\%$ rate determined from the inflow velocity.

Plain Language Summary When the solar wind impacts the Earth's magnetosphere, an explosive energy conversion process called magnetic reconnection opens the door for solar wind energy to enter the magnetosphere by interconnection of the magnetic fields of the solar wind and of Earth. In this process, magnetic energy is converted to charged-particle energy. Magnetic reconnection is fairly well understood at large scales and even down to the ion scale. However, the breaking and linking of field lines and the acceleration of electrons occur at much smaller scales, which are only recently being accessed by the NASA Magnetospheric Multiscale mission. This paper analyzes the speed at which electrons flow into and out of reconnection sites. The inflow speeds are crucial because they provide a measurement of the rate at which reconnection proceeds.

1. Introduction

During asymmetric reconnection, as occurs at the dayside magnetopause (MP), the flow stagnation point does not coincide with the X line, as it does with symmetric reconnection (Cassak & Shay, 2007; Priest et al., 2000). Since there is no transport of magnetic flux across the X line, and no net transport of mass flux across the stagnation point, the two points coincide only for equal mass densities, inflow velocities, and magnetic field strengths within each inflow region. The higher densities on the magnetosheath side and the

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higher magnetic field strength on the magnetosphere side displace the stagnation point toward the Earth from the X line.

At kinetic scales there are separate electron and ion stagnation points resulting from the larger gyroradii of ions (Cassak & Shay, 2009). Hesse et al. (2014) used particle-in-cell simulation and theory to examine energy conversion and electron distribution functions (DFs) within the asymmetric reconnection electron diffusion region (EDR). Their prediction of crescent-shaped electron distributions was confirmed with data from the NASA Magnetospheric Multiscale (MMS) mission by Burch, Torbert, et al. (2016). Pritchard et al. (2019) noted the occurrence of converging electron flows in a region of strong out-of-plane current and interpreted them as electron inflow velocities to the EDR. Inflow velocities are important because when scaled to the electron Alfvén speed (v_{Ae}), they provide the reconnection rate (R) (Karimabadi et al., 2013; Klimas, 2015). The use of v_{Ae} instead of v_{Ai} in this context is warranted by the fact that near the EDR the magnetic field is advected by the electrons (Cassak et al., 2005; Tsiklauri, 2008).

We report on a study of electron inflows associated with three MP reconnection events and one magnetosheath electron-only event. The inflow velocities (v_{eN}) range from 0.05 to $0.25v_{AeL}$ (with v_{AeL} = the asymmetric electron Alfvén speed based on the reconnecting magnetic field in the inflow regions). The outflow speeds (v_{eL}) have magnitudes 1–3 times v_{eN} , and the out-of-plane velocities (v_{eM}) are about 3–4 times v_{eN} . While the strong v_{eM} is well known as the main current contribution to $\mathbf{J}\cdot\mathbf{E}$ (ohmic dissipation), v_{eN} has not been analyzed before. We derive a reconnection rate of $0.04 \pm 25\%$ for one event using the reconnection electric field (E_M) as $R = E_M/(v_{AeL}B_L)$, with B_L the average L component of the magnetic field in the EDR, as compared to a value of $0.12 \pm 20\%$ for $R = v_{eN}/v_{AeL}$.

2. Data and Methods

MMS makes electron-scale measurements at four locations within or surrounding magnetic reconnection sites in the boundary regions of the magnetosphere. The measurements are summarized by Burch, Moore, et al. (2016), with details on plasma and electric and magnetic field data provided by Pollock et al. (2016) and Torbert et al. (2016). A major enabling factor for electron-scale reconnection studies is the unprecedented 30-ms time resolution of 3-D electron distributions by the Fast Plasma Investigation (FPI). For special studies, even faster measurements at 7.5 ms are possible because of the interleaved nature of the azimuthal sampling of FPI (Rager et al., 2018). These faster measurements are derived during ground processing and are used in parts of the current study.

3. Results

3.1. 15 April 2018 Event

MMS encountered the MP near 10 magnetic local time (MLT) at low negative geocentric solar magnetospheric (GSM) latitudes. Figure 1a shows the maximum shear model (Trattner et al., 2012) applied to this event with the MMS position noted by the square symbol. Figure 1b shows the S/C positions in boundary-normal coordinates.

Figures 1c–1h plot magnetic field, electron velocities, and average ion velocities for MMS1–4. There was a significant guide field (B_M) of about 0.4 times the magnetosphere reconnecting field. The method of Schwartz (1998) was used to estimate a reconnection structure velocity along N at -23 km/s. Following Denton et al. (2016), the structure velocity along L is estimated at -50 km/s based on $\langle v_{iL} \rangle$ as plotted in Figure 1f. The most prominent features in the electron velocity plots (Figures 1f–1h) are the peaks in v_{eM} in (g), which carry the out-of-plane current associated with the reconnection EDR. Using the peaks in v_{eM} as the most prominent markers, we note in Figure 1h that three of them are spanned by bipolar signatures in the inflow velocities v_{eN} . Only in MMS4 is the inflow velocity unipolar, which may indicate that it encountered just one side of the EDR. These bipolar signatures are most easily seen for MMS2 (red plots). Figure 1f shows that also overlapping the v_{eM} peaks is a single or binary v_{eL} jet. Such v_{eL} peaks can also be seen in Figure 3 of Pritchard et al. (2019).

Figure 2 shows results of a reconstruction of the reconnection magnetic field using the polynomial method of Denton et al. (2020). A very similar result was obtained using the method of Torbert et al. (2020). Figure 2a shows the magnetic field averaged over MMS1–4 with vertical dotted lines indicating times of the plots of magnetic field 1

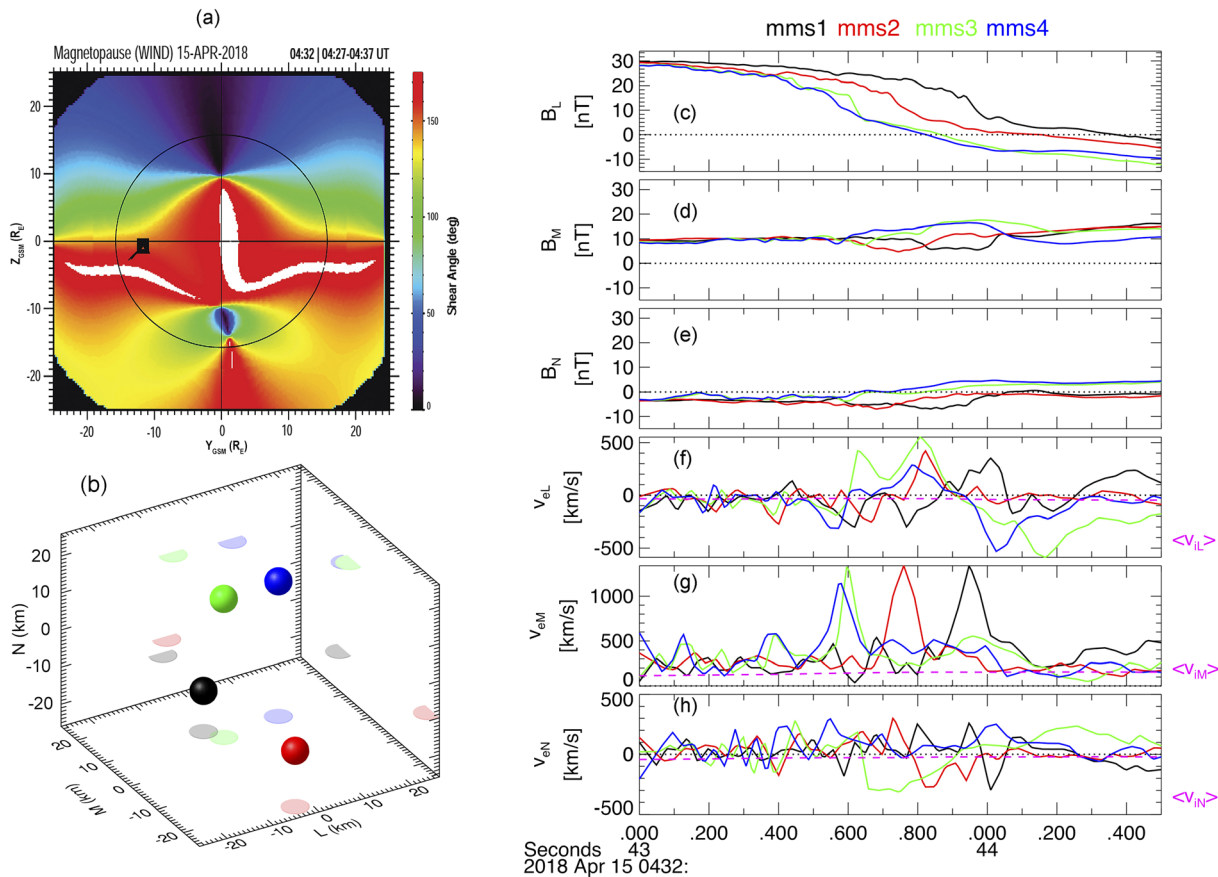


Figure 1. (a) Shear angle between magnetosheath and magnetosphere magnetic fields at 04:32 UT on 15 April 2018. Plot is in GSM coordinates as viewed from the Sun. MMS location is noted by the black square with the ion flow velocity indicated by the line attached to the square. Color coding shows a shear angle near 150° . (b) MMS1–4 locations at 04:32 UT with MMS1 in black, MMS2 in red, MMS3 in green, and MMS4 in blue. Plot is in boundary-normal coordinates (LMN) with GSE to LMN derived from the joint variance analysis (Genestreti et al., 2018; Mozer & Retino, 2007): $L = [0.21860787, 0.064063071, 0.97370762]$ GSE, $M = [-0.59993242, -0.77815588, 0.18588845]$ GSE, $N = [0.76960490, -0.62479544, -0.13167749]$ GSE. (c–e) Magnetic field and (f–h) electron velocity in LMN coordinates for the four MMS S/C. Also plotted with each v_e is the average v_i (in dashed magenta) across all four S/C.

ines in the LN plane (panels b–g). The vertical dotted lines also correspond to times of DF plots shown in Figures 3 and 4 with Figures 2d–2f corresponding to Figures 3j–3l and Figures 2b, 2c, and 2g corresponding to Figures 4j, 4k, and 4m.

3.1.1. Data From MMS2 and MMS3

Figures 3 and 4 show data for MMS2 and MMS3, respectively. MMS was inbound at the morning-side (~ 10 MLT) MP, but since the MP was moving inward at a higher velocity, the MP crossing was outbound as seen by the magnetic field plots in Figures 1–4. In these plots, the FPI electron velocities in panel (c) and pitch angle distributions (PADs) in panel (i) are derived from the 7.5-ms data. The E-field data in panel (d) are averaged to 7.5-ms resolution so that the $\mathbf{J}\cdot\mathbf{E}$ values in panels (g) and (h) are also at 7.5-ms resolution. As shown in Figures 3b, 3f, 4b, and 4f, the wave activity was predominantly electrostatic with frequency peaks between the electron cyclotron and plasma frequencies. Figures 3c and 4c both show a strong peak of electron velocity in the M direction; the width of these peaks at half maximum is $\sim 1.5 d_e$, where d_e is the electron inertial length of ~ 1.5 km.

3.1.2. Comparison of MMS2 and MMS3

We first focus on the primary energy conversion site, which for both S/C was within the electron stagnation region, although as seen in Figure 2b, MMS3 was near the magnetosphere separatrix. We define this location by the time of the $\mathbf{J}\cdot\mathbf{E}$ peaks (43.75 s for MMS2 and 43.61 s for MMS3).

Next, we note in Figures 3c and 4c that the bipolar patterns of v_{eN} and v_{eL} overlap the broad peak of v_{eM} ; but there is a difference in that for MMS2 these patterns are shifted only slightly to a later time, while this shift is

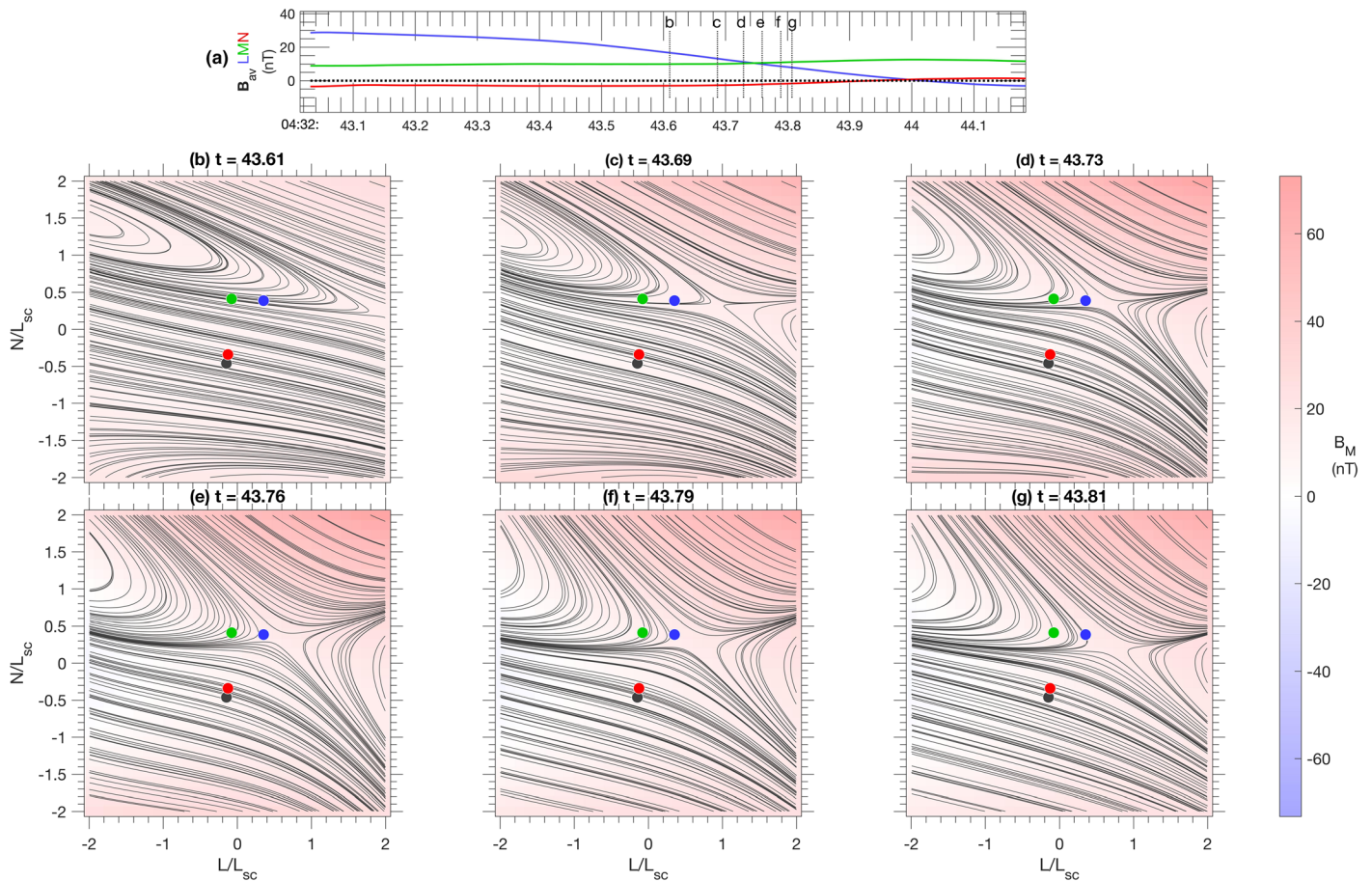


Figure 2. Reconstruction of the reconnection magnetic field for the 15 April 2018 event using the polynomial method of Denton et al. (2020). (a) Average LMN magnetic field over all four S/C. (b–g) L–N cuts through a 3–D model at times indicated by vertical dotted lines in (a) with the S/C positions indicated by the colors from Figure 1. The color bar indicates values of B_M . The vertical dotted lines are also time markers for electron distribution functions (DFs) plotted in Figures 3j–3l, 4j, 4k, and 4m. L_{sc} is the average inter-spacecraft distance.

more significant for MMS3. In fact, for MMS3 the inflow velocity $v_{eN} > 0$ occurs just past the v_{eM} peak. For MMS2 (Figure 3d) E_N and E_L stayed positive through the peak out-of-plane current while E_M maintained a small negative value, as expected for the reconnection E field. MMS3 (Figure 4d), on the other hand, detected a bipolar signature of higher values of E_M , which was accompanied by a slightly shifted bipolar signature of E_L . The maximum negative value of E_M during this bipolar trace coincided with the peak of $\mathbf{J} \cdot \mathbf{E}$. This type of pattern was reported before by Burch et al. (2018) at the boundary between open and closed field lines in an EDR, which is consistent with the MMS3 location shown in Figure 2b.

We now focus on the $\mathbf{J} \cdot \mathbf{E}$ plots in panels (g) and (h) of Figures 3 and 4. $\mathbf{J} \cdot \mathbf{E}$ is a scalar quantity, but it is instructive to plot separately the contributions from the L, M, and N components of \mathbf{J} and \mathbf{E} in panel (g). The plots in (g) and (h) are made in the rest frame of the reconnection structure so that $\mathbf{E}_{str} = \mathbf{E} + \mathbf{v}_{str} \times \mathbf{B}$. In Figure 3g it is notable that the green trace ($J_M E_M$) has a significant positive value through the broad v_{eM} peak shown in panel (c). There is a smaller, mostly positive blue trace ($J_L E_L$) but a strong negative red trace ($J_N E_N$). This negative $J_N E_N$ is caused by the positive E_N and the positive v_{eN} shown in panel (c), which is the inflow velocity. This alignment of E_N and v_{eN} agrees with the conclusion of Swisdak et al. (2018) that the oscillations in $\mathbf{J} \cdot \mathbf{E}$ arise from changes in the sign of v_{eN} .

In the total $\mathbf{J} \cdot \mathbf{E}$ plot in Figure 3h, the negative $J_N E_N$ is seen to cause net negative values. The $\mathbf{J} \cdot \mathbf{E}$ plots for MMS3 in Figures 4g and 4h show a much stronger and narrower $J_M E_M$ peak and only small negative values in the total $\mathbf{J} \cdot \mathbf{E}$ plot just before the peak. In this case the negative total $\mathbf{J} \cdot \mathbf{E}$ was caused by the positive values of E_M shown in panel (d) within the positive v_{eM} peak shown in panel (c).

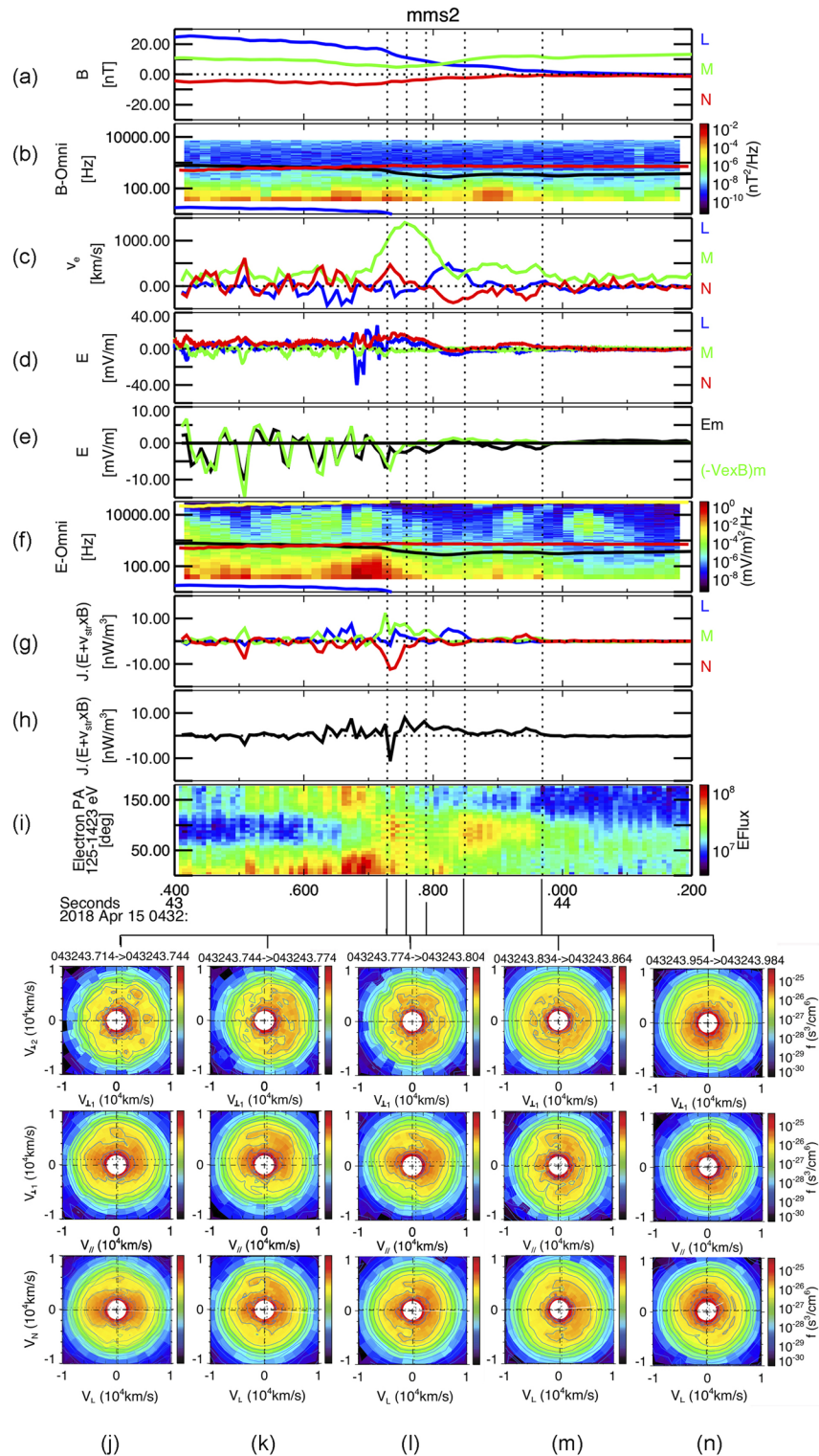


Figure 3. Data from MMS2 on 15 April 2018. (a) Magnetic field, (b) omnidirectional power spectral density (PSD) of B, (c) electron velocity in spacecraft frame, (d) electric field in spacecraft frame, (e) comparison of E_M and $(-v_e \times B)_M$, (f) omnidirectional PSD of E, (g) contributions to $\mathbf{J} \cdot \mathbf{E}$ in the rest frame of the reconnection structure, (h) total $\mathbf{J} \cdot \mathbf{E}$, in the structure frame, (i) electron pitch angle distributions (PADs) for 125–1423 eV with EFlux as $\text{eV cm}^{-2} \text{sr}^{-1} \text{eV}^{-1}$, and (j–n) electron distribution function (DF) top: in plane perpendicular to \mathbf{B} with $v_{\perp 1}$ along $\mathbf{E} \times \mathbf{B}$ and $v_{\perp 2}$ along \mathbf{E} ; middle: in plane containing \mathbf{B} and $v_{\perp 1}$; bottom: v_L versus v_N . The yellow, red, black, and blue traces in the wave spectrograms plot F_{pe} (electron plasma), F_{pi} (ion plasma), F_{ce} (electron cyclotron), and F_{lh} (lower hybrid) frequencies.

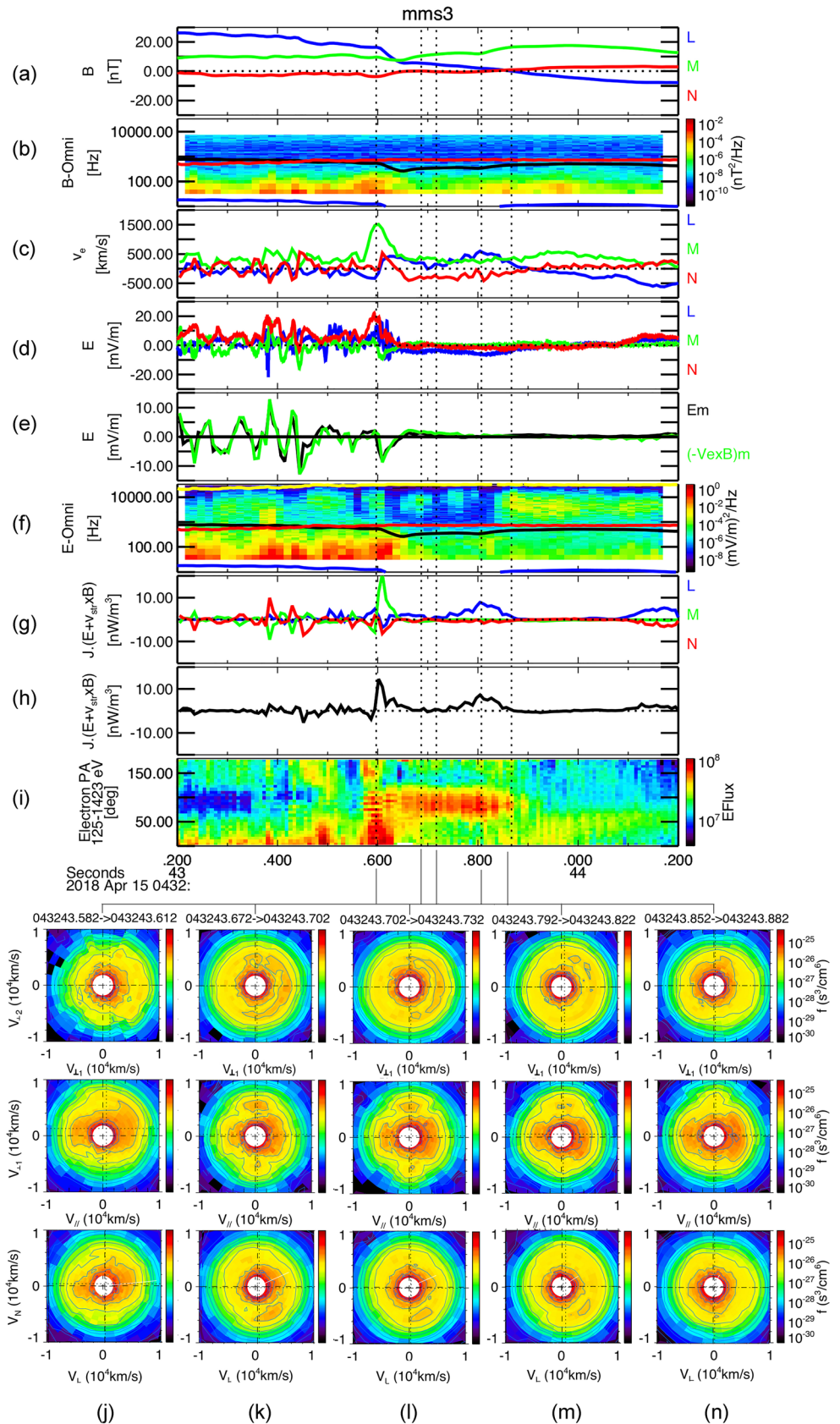


Figure 4. Same as Figure 3 except for MMS3.

Looking now at the region between the **J-E** peak and the B_L reversal, Figure 3g for MMS2 shows a peak in $J_L E_L$ associated with the outflow v_{eL} in panel (c) and a peak in $J_N E_N$ associated with a localized increase in the flow along $-N$ noted above. The MMS3 data in Figure 4 are different in that the only secondary **J-E** peak occurred near the B_L reversal where there is a peak in v_{eL} within a broader region of $v_{eL} > 0$ and $E_L < 0$. The peak in $v_{eL} > 0$ just before the B_L reversal, which is not seen by MMS2, is interpreted as electron flow toward the X line in the magnetosheath boundary layer as shown by the simulations of Swisdak et al. (2018) just to the $-L$ side of the X line. Moving to the left in Figure 4c, we note an enhancement in $-v_{eN}$ (earthward flow) and a decrease in v_{eL} , suggesting a clockwise turning of the flow from the X line to the Earth, which is complete between the second and third vertical dotted lines where v_{eL} drops to zero.

3.2. Electron Flow Observations on 15 April 2018, 14 December 2015, 9 December 2016, and 29 December 2016

Figure 5 shows electron flow velocities for four reconnection events. The two vertical dotted lines mark the peak inflow velocities, which we identify as the edges of the EDR. Figures 5a–5e show the event on 15 April 2018, which was presented in detail in section 3.1. The inflow speeds approached $0.14v_{AeL}$, and the peak outflow speed was only slightly larger than the inflow speeds, but with the out-of-plane speed reaching $\sim 0.7v_{AeL}$, following the same pattern as the other events. Also shown for this event are the 8,175/s electric field components in the X-line frame and correlations between 10-point average E_M and E_N (panel g) and E_L (panel h), which are discussed in section 3.3. The second event, on 14 December 2015 (panels l–p), has been the subject of previous publications by Chen et al. (2017) and Graham et al. (2017) among others. The inflow speed reached $\sim 0.12v_{AeL}$, implying a reconnection rate near 0.12 while the outflow speeds (panel e) were about twice the inflow speeds. Similar patterns are seen for the event on 9 December 2016 (panels q–u), which is the magnetosheath electron-only reconnection event reported by Phan et al. (2018). This event is included because, although not discussed by Phan et al., electron inflows in the form of bipolar v_{eN} were present and provide another valuable comparison. In this case the inflow speeds reached $\sim 0.25v_{AeL}$, implying a reconnection rate of ~ 0.25 . MMS2 measured a unipolar outflow speed (panel u) of up to $0.4v_{AeL}$, while the highest speeds were observed in the out-of-plane flow (panel t), which reached nearly $0.8v_{AeL}$.

Comparable patterns are seen for the event on 29 December 2016 (Figures 5v–5z), which occurred in the electron stagnation region. For this event, which was published by Pritchard et al. (2019), the inflow speeds reached $\sim 0.05v_{AeL}$, indicating a reconnection rate of ~ 0.05 .

3.3. Electric Field Data and Error Analysis

The tilt of the X-line structure in the L-N plane in the Figure 2 reconstruction could be due to a combination of a sub-optimal LMN transform and inaccuracies in the reconstruction code. As shown in Figures 3e and 4e, there is very good equivalence between E_M and $(-v_e \times B)_M$ except, as expected, when MHD is violated in dissipation regions as identified by $\mathbf{J} \cdot \mathbf{E}_{\text{str}}$. Thus, we conclude as have others (e.g., Torbert et al., 2017) that the measurements of v_e are very accurate and that errors in the LMN transformation represent the primary measurable source of error in the reconnection rate. Following Genestreti et al. (2018), we can assess the accuracy of the LMN transform by comparing the values of E_M , which we assume are uniform within the EDR, with the larger values of E_L and E_N . Plots of these comparisons are shown in Figures 5g and 5h. The slopes of the fit lines in these plots give the tilt angle of the two axes in radians since for small angles $\cos(\text{tilt}) \sim \text{tilt}$. From Figure 5g the tilt angle in the N,M plane between the ideal and actual M axes is $\sim 10.7^\circ$, while Figure 5h shows the tilt in the L,M plane to be $\sim 10.1^\circ$. Since our LMN transform requires orthogonality, we conclude that a similar error occurs in the L,N plane.

We have applied a rotation to the electric field data using the matrix shown in the Figure 5 caption. This rotation minimizes the contamination of E_M values by E_N and E_L . The mean value of E_M is -1.164 mV/m with a standard deviation of 0.376. The normalized reconnection rate derived from E_M is $\langle E_M \rangle / (v_{AeL} \langle B_L \rangle) \sim 0.04 \pm 25\%$. In order to estimate the error in the reconnection rate derived from v_{eN}/v_{AeL} , we applied the same correction of the LMN transform to the electron velocity data and found a reconnection rate of 0.12, which is $\sim 16\%$ lower than shown at the vertical dashed lines in Figure 5c for the original LMN transform. Thus, we conclude that the relative error in the reconnection rate derived from the electron inflow velocity is about $\pm 20\%$, which is similar to the error in E_M .

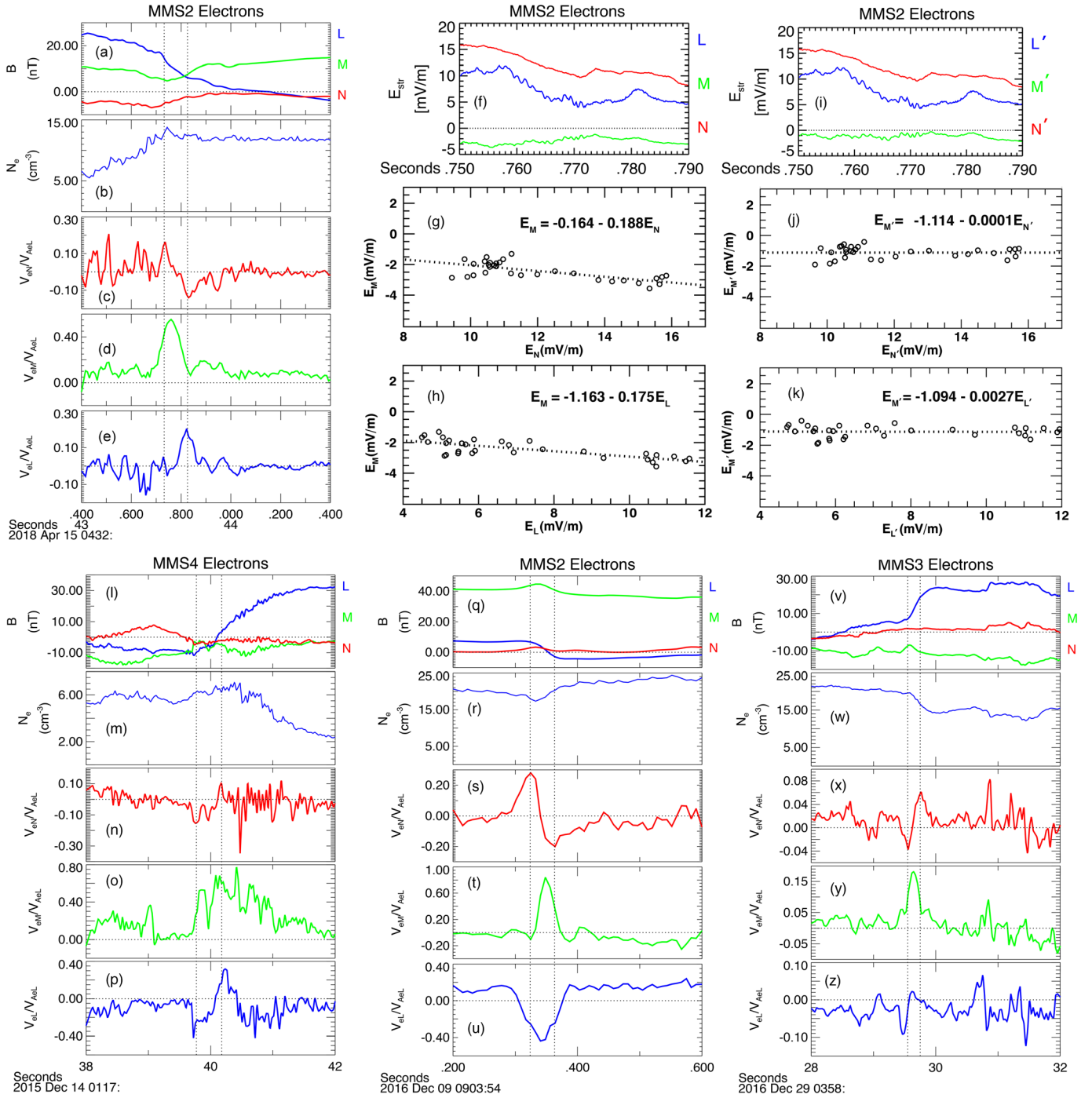


Figure 5. Electron velocities for reconnection events on 15 April 2018, 14 December 2015, 9 December 2016, and 29 December 2016. For each event there are plots of (a, l, q, v) B_{LMN} , (b, m, r, w) N_e , (c, n, s, x) v_{eN}/v_{AeL} , (d, o, t, y) v_{eM}/v_{AeL} , and (e, p, u, z) v_{eL}/v_{AeL} , where v_{AeL} is the electron Alfvén speed given by $v_{AeL}^2 = B_{L1}B_{L2}(B_{L1} + B_{L2})/(\mu_0\rho_1B_{L2} + \mu_0\rho_2B_{L1})$ with subscripts 1 and 2 denoting the first and second vertical dotted lines, respectively (see Cassak & Shay, 2007). Values of v_{AeL} for the four events are 2,497, 2,177, 946, and 2,488 km/s, respectively. Also, (f) electric field in X-line rest frame, (g) correlation between 10-point average E_N and E_M , (h) correlation between 10-point average E_L and E_M , (i) E field with rotation correction, (j) E_N , E_M correlation with rotation correction, and (k) E_L , E_M correlation with rotation correction. Rotation correction matrix given by N' : $[\sin(p)\sin(ts), \sin(p)\cos(ts), \cos(p)]$, M' : $[\cos(p)\sin(ts), \cos(p)\cos(ts), -\sin(p)]$, L' : $[\cos(ts), -\sin(ts), 0]$, where $p = -0.188$ rad (from panel g), and $ts = 0.175$ rad (from panel h).

4. Conclusions

We have presented electron velocities within an MP reconnection site on 15 April 2018. The observations were made at all four spacecraft with detailed data shown for MMS2, which was well within the electron stagnation region, and MMS3, which was near the magnetosphere separatrix. In both cases electron outflow jets (along L and/or $-L$) and inflow velocities (along $\pm N$) were observed to span the region of highest out-of-plane velocity and $\mathbf{J} \cdot \mathbf{E}$. These observations are compared to previously published events on 14 December 2015, 9 December 2016, and 29 December 2016 and found to have similar values in that the inflow velocities were in the range from 0.05 to $0.25v_{eA}$, which provide normalized reconnection rates. We used electric field data to estimate the inaccuracy of the LMN transform for the 15 April 2018 event to be $\sim 10.5^\circ$. We corrected the transform errors by applying an additional matrix rotation and obtained a normalized reconnection rate of $0.04 \pm 25\%$. Applying the same LMN correction to the electron velocity data yielded a reconnection rate about 16% below that shown in Figure 5c or ~ 0.12 with an estimated error of $\pm 20\%$. The difference between the reconnection rates derived from E_M and v_{eN} is not understood and will be investigated for more events. One possible explanation is the existence of v_{eN} components parallel to \mathbf{B} , which would not contribute to the advection of the magnetic field and so would reduce the reconnection rate derived from the inflow velocity. The mean angle between v_{eN} and \mathbf{B} in Figure 5c is $87.6^\circ \pm 19\%$, which places this potential error within the 20% range derived from the LMN transform uncertainty.

The outflow velocities along $\pm L$ ranged up to about three times the inflow velocities. This ratio is similar to the theoretical prediction of the aspect ratio of the inner EDR based on electron trapping length in a field reversal (Hesse et al., 1999). A similar result (aspect ratio of inner EDR ~ 4) was obtained with MMS for a tail reconnection event by Nakamura et al. (2019) who found $v_{eL} \sim 0.1v_{eA}$ in the outflow region. It is predicted that such sub-Alfvénic outflow in the inner EDR will accelerate to v_{eL} as the electrons move toward the exhaust region as shown by the simulations of Shay et al. (2007) and Drake et al. (2008).

Similar investigations are conducted in the laboratory with MRX: Yamada et al. (2018) observed the high out-of-plane velocity (v_{eM}) at the stagnation point of asymmetric reconnection; Ren et al. (2008) observed outflow velocities at $0.11v_{eA}$, which are consistent with our measurements.

Data Availability Statement

The 15 April 2018 reconnection event was discovered in a database created for the International Space Science Institute Team 442, "Study of the physical processes in magnetopause and magnetosheath current sheets using a large MMS database." The entire MMS data set is available online (at <https://lasp.colorado.edu/mms/sdc/public/links/>). Fully calibrated data are placed online at this site within 30 days of their transmission to the MMS Science Operations Center. The data are archived in the NASA Common Data Format (CDF) and so can be plotted using a number of different data display software packages that can use CDF files. A very comprehensive system called the Space Physics Environment Data Analysis System (SPEDAS) is available by downloading (http://themis.ssl.berkeley.edu/socware/bleeding_edge/) and selecting (spdsw_latest.zip). Training sessions on the use of SPEDAS are held on a regular basis at space physics-related scientific meetings. All of the data plots in this paper were generated with SPEDAS software applied to the publicly available MMS database, so they can readily be duplicated.

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