

Ocean-atmosphere coupled Pacific Decadal variability simulated by a climate model

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Abstract

25
26 Currently, the mechanisms for Pacific Decadal Oscillation (PDO) are still
27 disputed, and in particular the atmosphere response to the ocean in the mid-latitude
28 remains a key uncertainty. In this study, we investigate two potential feedbacks – a
29 local positive and a delayed negative – for the PDO based on a long-term control
30 simulation using the ECHAM5/MPI-OM coupled model, which is selected because of
31 reproduces well the variability of PDO. The positive feedback is as follows. In the
32 PDO positive phase, the meridional sea surface temperature (SST) gradient is
33 intensified and this strengthens the lower level atmospheric baroclinicity in the
34 mid-latitudes, leading to the enhancement of Aleutian low and zonal wind. These
35 atmospheric changes reinforce the meridional SST temperature gradient through the
36 divergence of ocean surface currents. The increased heat flux loss over the
37 anomalously warm water and decreased heat flux loss over the anomalously cold
38 water in turn reinforce the lower atmospheric meridional temperature gradient,
39 baroclinicity and atmospheric circulation anomalies, forming a local positive feedback
40 for the PDO. The delayed negative feedback arises, because the intensified meridional
41 SST gradient also generates an anticyclonic wind stress in the central North Pacific,
42 warming the upper ocean by Ekman convergence. The warm upper ocean anomalies
43 then propagate westward and are transported to the mid-latitudes in the western North
44 Pacific by the western boundary current. This finally reduces the meridional SST
45 gradient, 18 years after the peak PDO phase. **These results demonstrate the significant**
46 **contributions of the meridional SST gradient to the PDO's evolution.**

47 **Key words:** PDO, barotropic structure, feedback, meridional SST gradient,

48 ECHAM5/MPI-OM

49 **1. Introduction**

50 The Pacific Decadal Oscillation (PDO) is a pattern of mid-latitude climate
51 variability on decadal to inter-decadal time scales over the North Pacific (Mantua et al.
52 1997). It is commonly defined as the leading empirical orthogonal function (EOF) of
53 North Pacific (20°–60°N) SST monthly anomalies, which are departures from the
54 climatological annual cycle after regressing out the global mean SST changes
55 (Newman et al. 2016). The associated principal component time series is usually
56 regarded as the PDO index. The PDO affects the climate system in two ways. First,
57 PDO perturbations are superimposed on long-term climate trends, thereby affecting
58 decadal variations in climate both locally and globally. For example, PDO influences
59 decadal climate variability over the North Pacific Ocean, North America and East
60 China (Latif and Barnett 1994, 1996; Mantua et al. 1997; Yang et al. 2017) and likely
61 contributed to the recent hiatus in global warming (Trenberth and Fasullo 2013).
62 Second, the PDO provides the background state for interannual variability and may
63 modulate it. For instance, the PDO exerts a modulating effect on ENSO (Lin et al.
64 2018; Vimont et al. 2001; Vimont et al. 2003) and its teleconnections (Gershunov and
65 Barnett 1998; Wang et al. 2008; Wang and Liu 2015). Thus, the PDO has important
66 impacts on social and economic development globally, and in particular in North
67 America and East Asia.

68 Although numerous studies on the PDO and its mechanisms have been
69 conducted since 1990, our understanding of the processes driving it remains limited,
70 especially the existence of an atmospheric response to mid-latitude SST and its role in

71 the PDO (Newman et al. 2016). The atmospheric signature of the PDO is an
72 equivalent barotropic vertical structure, with low (high) pressure above cold (warm)
73 water (Cayan 1992a, b). The influence of the atmosphere on the mid-latitude Pacific
74 Ocean is well understood: SST is driven directly via thermodynamic interaction and
75 indirectly via changes in ocean circulation (Alexander et al. 2002; Cayan 1992a, b;
76 Latif and Keenlyside 2011; Liu and Alexander 2007). Early studies considered the
77 impact of mid-latitude ocean on the atmosphere to be weak in comparison to the
78 strong intrinsic variability of mid-latitude atmosphere (Kushnir et al. 2002).
79 Furthermore, models disagree on the atmospheric responses to mid-latitude SST
80 anomalies. Some show a barotropic response, with low pressure over cold water, or
81 with high pressure over cold water (Kushnir and Lau 1992) and others suggest a
82 baroclinic one (Hoskins and Karoly 1981). These different results from models are
83 related to differences in the atmospheric heating profiles caused by the SST anomalies
84 (Frankignoul 1985). However, recent studies based on higher resolution observations
85 and models suggest that the atmosphere does respond to the SST anomalies in western
86 boundary current regions, particularly the subarctic frontal zone (Ma et al. 2015; Ma
87 et al. 2016; Nakamura et al. 2004; O'Reilly and Czaja 2015; Small et al. 2008). In this
88 region, persistent SST anomalies in winter tend to change the state of atmosphere
89 through heat and moisture fluxes (Taguchi et al. 2012), and these atmosphere
90 anomalies in turn affect the ocean state (Qiu and Chen 2005; Qiu et al. 2007).
91 Although possibly weaker than observed, these feedbacks might be active also in
92 lower resolution climate models (Latif 2006; Latif and Barnett 1994, 1996; Park and

93 Latif 2010; Zhang and Delworth 2015). Thus, this atmospheric response to
94 mid-latitude SST anomalies can be important in decadal climate variability (Fang and
95 Yang 2016; Kwon and Deser 2007; Zhong et al. 2008; Zhou et al. 2015), and
96 understanding it is critical for the simulation and prediction of PDO.

97 The mechanisms raised to driving the PDO in previous studies could be broadly
98 divided into three categories. First, the PDO is driven by the stochastic forcing of the
99 atmosphere. From this perspective, high frequency weather fluctuations excite low
100 frequency ocean variations (e.g., Frankignoul et al. 1997; Jin 1997; Latif and
101 Keenlyside 2011). While the stochastic forcing does not depend on an atmospheric
102 response to the developing ocean anomalies, it can be important in sustaining a
103 damped coupled mode (Latif and Keenlyside 2011). Furthermore, an atmospheric
104 response to mid-latitude SST anomalies could enhance the decadal variance, or lead
105 to a preferred time scale, if the response is a positive or delayed negative, respectively,
106 feedback on the low-frequency variability (Barsugli and Battisti 1998; Frankignoul
107 1985; Lee et al. 2008; Sura et al. 2006). Second, the PDO is driven by teleconnection
108 from the tropics (Newman et al. 2003; Shakun and Shaman 2009). In this case,
109 tropical variability, such as ENSO or the decadal variability of ENSO, drives
110 variability in the North Pacific by the atmosphere (e.g., Alexander et al. 2002;
111 Trenberth et al. 1998) and/or ocean (e.g., Gu and Philander 1997; Kleeman et al.
112 1999). Besides, two-way interactions between the tropical and North Pacific may
113 impact low-frequency variability in both domains (e.g., Newman 2007; Solomon et al.
114 2003). Third, the PDO is driven by mid-latitude air–sea interactions. In this viewpoint,

115 the unstable ocean–atmosphere interactions in the mid-latitude would lead to the PDO
116 (Fang and Yang 2011, 2016; Latif and Barnett 1994, 1996; Zhang and Delworth 2015).
117 Besides, compared to the limited predictive skill resulting from stochastic forcing,
118 predictability may be potentially enhanced by taking into account the air-sea
119 interaction (Barnett et al. 1999; Pierce 2001). Thus, the PDO can be regarded as a
120 result of a combination of different physical processes, but the relative importance of
121 these, especially the third, remains uncertain (Newman et al. 2016). In this study, a
122 coupled general circulation model (CGCM) developed at the Max Planck Institute
123 that reproduces well the key features of the PDO is used to investigate the structure
124 and possible feedback mechanisms of the PDO.

125 The remainder of the paper is organized as follows. The model and data used are
126 described in Section 2. In Section 3, the typical structure of the wintertime PDO in
127 CGCM is presented. Feedbacks in CGCM to maintain or reverse the phase of the
128 PDO are examined in Sections 4.1 and 4.2, respectively. The final section is the
129 conclusions and discussion.

130

131 **2. Methodology**

132 *2.1. Model*

133 The current analysis employs a coupled atmosphere-ocean general circulation
134 model of the Max Planck Institute (ECHAM5/MPI-OM) (Jungclaus et al. 2006) to
135 explore the possible mechanism for the PDO. The fifth-generation atmospheric
136 general circulation model (ECHAM5) has been developed from the European Center

137 for Medium range Weather Forecasting operational forecast model cycle 36 (1989)
138 and a comprehensive parameterization package developed at Hamburg. It is run at
139 T63 resolution ($1.875^\circ \times 1.875^\circ$) with 31 vertical (hybrid) levels. The Max–Planck–
140 Institute ocean model (MPI-OM) has been derived from the Hamburg Ocean
141 Primitive Equation model, which simulates the primitive equations for a hydrostatic
142 Boussinesq fluid with a free surface and includes an embedded sea ice model. It has
143 1.5° average horizontal grid spacing with 40 unevenly spaced vertical levels from 6 m
144 to 5720 m. More technical details of ECHAM5 and MPI-OM can be found in
145 Roeckner et al. (2003) and Marsland et al. (2003). The atmosphere and ocean are
146 coupled once a day via the Ocean–Atmosphere–Sea Ice–Soil (OASIS) coupler
147 (Valcke et al. 2003). The ocean passes SST, sea ice concentration, sea ice thickness,
148 snow depth, and ocean surface velocities to the atmosphere, whereas the atmosphere
149 passes surface wind stress, fresh water flux, and heat flux to the ocean. No flux
150 adjustment is applied. This study use monthly data from a pre-industrial control
151 experiment of ECHAM5/MPI-OM that excludes the effect of human activities by
152 imposing non-evolving pre-industrial conditions. Thus, here we consider the PDO as
153 a natural pattern of climate variability resulting from dynamics internal to the climate
154 system. In this study, we conduct a 600–year control simulation (Zheng 2014), and
155 analyze the last 150 years of the result; the PDO in this period shows very similar
156 characteristics to the PDO of the entire run.

157

158 **2.2. Data**

159 We use monthly mean observations from 1854 to 2014, including sea surface
160 temperature (SST) and atmospheric circulation data, to evaluate the performance of
161 ECHAM5/MPI-OM. The SST data are the Extended Reconstructed Sea Surface
162 Temperature version 3b (ERSST V3b) (Smith et al. 2008): a global SST analysis with
163 a 2° horizontal resolution and derived from the International Comprehensive Ocean–
164 Atmosphere Dataset with missing data filled in by statistical methods. The
165 atmospheric circulation data are the Twentieth Century Atmospheric Reanalysis
166 version 2c (20CRv2c) (Compo et al. 2011), which assimilates only surface pressure
167 reports and uses observed monthly sea-surface temperature and sea-ice distributions
168 as boundary conditions. It provides an estimate of the global tropospheric variability
169 with a 2° spatial resolution.

170 In this study, winter (DJF) refers to the months of December, January, and
171 February. Interannual anomalies are obtained by removing the mean seasonal cycle
172 from the original fields. Furthermore, the atmospheric baroclinicity is represented by
173 the Eady growth rate, which is calculated according to the relation

$$174 \quad \sigma_{BI} = \frac{0.31f \frac{\partial u}{\partial z}}{N}$$
$$175 \quad N^2 = \frac{g}{\bar{\theta}} \frac{d\bar{\theta}}{dz}$$

176 where f is the Coriolis parameter, u is the zonal wind, N^2 is the buoyancy
177 frequency or the Brunt–Vaisala frequency, g is the gravitational acceleration, and $\bar{\theta}$
178 is the potential temperature of the background state.

179

180 **3. Structure of the PDO**

181 ECHAM5/MPI-OM realistically simulates the PDO SST structure and timescale.
182 As in the observations, the model simulated PDO shows SST variations with negative
183 anomalies over the central and western mid-latitude North Pacific, surrounded by
184 positive anomalies (Fig. 1a,b). Besides, the maximum zone of negative climatological
185 meridional SST gradient nearly overlays in the western cooling region during the
186 PDO positive phase, implying the SST gradient increasing south of cooling center
187 while decreasing north of that. In the model, however, the negative anomalies are
188 located further to the west compared to observations and the positive anomalies
189 further to the North. Furthermore, the coupled model captures the main periodicities
190 of the observed PDO (Figs. 1c-f).

191 The atmospheric and oceanic structure of the PDO in DJF is investigated by
192 composite analysis in both observations and simulations, considering the difference
193 between periods when the index is above and below one standard deviation. During
194 the positive PDO phase, SST is anomalously cold in the western-to-central North
195 Pacific but warm off the west coast of North America and in low latitudes (Fig. 2).
196 This pattern leads to an southward shift of the SST gradient in the western Pacific
197 around 40°N. The SLP becomes low over the North Pacific, reflecting an enhanced
198 Aleutian low. Over the cooling SST center, the air temperature is anomalously cold in
199 lower layers of the troposphere but warm in upper layers. In contrast to the variation
200 of air temperature, the geopotential height decreases in the whole troposphere. This
201 observed relationship between the atmosphere and the ocean also exists in our

202 simulations, although the simulated structure locates further west compared to the
203 observation. In summary, north-east of the cold SST anomaly, the geopotential height
204 anomaly displays an equivalent barotropic low from the surface to 100 hPa. This kind
205 of vertical structure is dubbed an equivalent barotropic cold/trough structure, and
206 characterizes the configuration of the mid-latitude North Pacific ocean–atmosphere
207 anomalies on the decadal time scale.

208 It should be noticed that the above-mentioned vertical structure differs from that
209 in the tropics that is primarily thermal-driven (Gill 1980; Matsuno 1966). In the
210 thermal-driven situation, when the underlying temperature becomes cold, a baroclinic
211 response in the pressure field occurs, with increases in geopotential height at lower
212 layers and decreases in upper layers. Therefore, the feedback in the mid-latitude North
213 Pacific should also be considered differently from that in the tropics (Hoskins and
214 Karoly 1981). In the following sections, we focus on a positive feedback that
215 enhances this structure, and a delayed negative feedback that causes the transition
216 between phases.

217

218 **4. Mechanism for the PDO in ECHAM5/MPI-OM**

219 ***4.1. Positive feedback of the PDO***

220 Given the typical time scale of the mid-latitude upper ocean, the mid-latitude
221 ocean–atmosphere interaction is often considered as a potential source of
222 decadal-to-interdecadal climate variability (Mantua et al. 1997; Trenberth 1990).
223 Therefore, we focus on the region within 20°N–60°N, and 150°E –150°W to

224 investigate the local ocean–atmosphere feedback, through linear regression of DJF
225 anomalies on the wintertime PDO index. This region is indicated in Fig. 2 by a yellow
226 box.

227 Climatologically, the relationship between the atmospheric transient eddy
228 activities and the mean flow were previously identified with zonal averages over
229 whole Northern Hemisphere (Holopainen et al. 1982; Lau and Holopainen 1984).
230 Meanwhile, the transient eddy activities build a bridge of the atmospheric response to
231 midlatitude SST anomalies. Thus, we use zonal averages over the midlatitude North
232 Pacific (the yellow one in Fig. 2) to identify the atmospheric response to midlatitude
233 SST anomalies. The zonal average meridional gradient of SST is negative in this
234 region. Two maximum values exist at around 26°N and 42°N that correspond to the
235 subtropical and subarctic oceanic frontal zones, respectively (Fig. 3a). The subtropical
236 oceanic frontal zone is associated with wind–driven submesoscale subduction
237 (Hosegood et al. 2013), whereas the subarctic zone is caused by the confluence of
238 warm Kuroshio and cold Oyashio (Talley et al. 2011). During the PDO positive phase,
239 the tripole SST anomalies with cooling in the mid-latitude western North Pacific and
240 warming to the north and south of it (Fig. 2) causes the meridional gradient of SST to
241 intensify to the south of 40°N and weaken north of it (Fig. 3a), implying that the
242 subarctic oceanic front zone moves southward. As the language of ocean–atmosphere
243 communication, the surface heat flux plays an important role in the information
244 transfer between the ocean and atmosphere. And differential air-sea heat exchanges
245 across an oceanic frontal zone maintain lower level atmospheric baroclinicity (e.g.,

246 Kuwano-Yoshida and Minobe 2017; Nakamura et al. 2004; Taguchi et al. 2012). The
247 ocean is cooler than the atmosphere, and obtains fresh water in the north of 40°N;
248 while the ocean is warmer than the atmosphere, and loses fresh water in the south.
249 Thus, the sensible heat flux as well as the latent heat flux would be transferred to the
250 ocean in the north of 40°N while to the atmosphere in the south of 40°N. The
251 anomalous net surface heat flux acts as a diabatic heating source of the lower
252 troposphere to the south of 40°N and a heating sink to the north, and these are
253 determined by the latent and sensible heat flux. It is worth noting that these characters
254 in Fig. 3a are consistent to those in the KOE region, as the air-sea flux are strongest in
255 the KOE region (Kelly et al. 2010; Kwon et al. 2010; Qiu et al. 2007; Taguchi et al.
256 2012). Thereby the ocean enhances the meridional air temperature gradient in low
257 levels, with a maximum at 600 hPa (Fig. 3b). According to the thermal wind equation,
258 the vertical shear of zonal wind is proportional to the meridional gradient of
259 temperature. Therefore, the change in the Eady growth rate (Fig. 3c) is similar to that
260 of the meridional gradient of temperature, indicating that the ocean maintains the
261 anomalous mid-to-low tropospheric baroclinicity (Fig. 3c). This favours an increase in
262 atmospheric transient eddies, and in the poleward flux of heat and momentum in the
263 atmosphere, that alters the large-scale atmospheric circulation (Brayshaw et al. 2008;
264 Holton 2004; Hoskins and Valdes 1990; Nakamura et al. 2008). Besides these,
265 previous studies highlight the importance of diabatic heating on this issue: The latent
266 heating anomalies to the south likely explain the mid-tropospheric heating that has
267 been linked to synoptic scale variations (Papritz and Spengler 2015). And there is

268 evidence that such diabatic heating associated with anomalies in SST fronts
 269 influences extreme synoptic scale variability and in turn the Aleutian low in late
 270 winter (Kuwano-Yoshida and Minobe 2017).

271 Thus, the influence of mid-latitude ocean anomaly on the atmosphere may be
 272 through not only diabatic heating but also the transient eddies. To assess these
 273 processes quantitatively, an analysis based on the quasi-geostrophic potential vorticity
 274 (QGPV) equation is performed. The QGPV equation (Fang and Yang 2016; Taguchi et
 275 al. 2012) is described as follows:

$$276 \quad \left(\frac{\partial}{\partial t} + \bar{\vec{V}}_h \cdot \nabla \right) \left[\frac{1}{f} \nabla^2 \bar{\phi} + f + \frac{\partial}{\partial p} \left(\frac{f}{\sigma_1} \frac{\partial \bar{\phi}}{\partial p} \right) \right] = -f \frac{\partial}{\partial p} \left(\frac{\alpha}{\sigma_1} \frac{\bar{Q}_d}{\bar{T}} \right) - f \frac{\partial}{\partial p} \left(\frac{\alpha}{\sigma_1} \frac{\bar{Q}_{eddy}}{\bar{T}} \right) +$$

$$277 \quad \bar{F}_{eddy},$$

278 where overbars indicates the seasonal mean. The variables ϕ , \vec{V}_h , f , α , σ_1 and T
 279 indicate the geopotential height, the geostrophic wind, the Coriolis parameter, the
 280 Specific volume, the static stability parameter and the temperature, respectively. \bar{Q}_d
 281 is the seasonal-mean diabatic heating derived from the residual of the thermodynamic
 282 equation (Yanai and Tomita 1998). \bar{Q}_{eddy} is the seasonal-mean transient eddy
 283 heating expressed as,

$$284 \quad Q_{eddy} = -\nabla \cdot \overline{\vec{V}'_h T'} - \frac{\partial \overline{\omega' T'}}{\partial p} + \frac{R}{C_p P} \overline{\omega' T'},$$

285 where the prime indicates the departure from the seasonal mean. The variables ω , R
 286 and C_p indicate the vertical velocity, the gas constant for dry air and the specific heat
 287 of dry air at constant pressure. \bar{F}_{eddy} is the transient eddy vorticity forcing written
 288 as,

$$289 \quad \bar{F}_{eddy} = -\nabla \cdot \overline{\vec{V}'_h \zeta'},$$

290 where variable ζ indicates the vorticity.

291 As Figure 4 shows, the diabatic heating is centered below the middle troposphere,
292 with a maximum at 850 hPa (Fig. 4a), unlike that induced by convection in the tropics.
293 According to the QGPV equation, this shallow diabatic heating tends to decrease the
294 geopotential height below 850hPa and increase above. Thus, the Aleutian low is
295 deepened near the surface via the diabatic heating. And the thermal transport by the
296 transient eddy activities converges north of 40°N with a maximum at 300 hPa (Fig.
297 4b). This transient eddy heating centered at the mid-troposphere can also reduce the
298 geopotential height in a similar manner as the diabatic heating. Thus, both diabatic
299 heating and transient eddy heating create a baroclinic atmospheric response, which
300 suggest to deepen the Aleutian low in the lower-to-middle troposphere. Furthermore,
301 transient eddies also transport vorticity flux in the upper troposphere. The
302 convergence of vorticity flux by transient eddy activities produces a maximum at
303 300hPa (Fig. 4c), which generates a barotropic response and reduces the geopotential
304 height in the upper troposphere. Consequently, as a joint result of thermal and eddy
305 driving, the Aleutian low is enhanced in the whole troposphere (Fig. 4d).

306 Accompanied with the deepened Aleutian low, the surface zonal wind is
307 enhanced as shown in Fig. 5a, especially the westerly wind to the north of negative
308 SST anomalies, resulting in changes in ocean circulation. To investigate the influence
309 of ocean circulation changes in SST, the SST tendency induced by advctions, namely
310 $-(\vec{V} \cdot \nabla T)' = -\left(u \frac{\partial T}{\partial x} + v \frac{\partial T}{\partial y} + w \frac{\partial T}{\partial z}\right)'$, is displayed in Fig. 5b. The anomalous
311 advection of SST can cause a large-scale SST cooling in the North Pacific, which

312 covers the cooling center during the PDO positive phase. Therefore, through the
313 ocean dynamical processes forced by the increased surface westerly, the original cold
314 SST anomaly will be sustained or even intensified, leading to the maintenance or
315 enhancement of the SST gradient anomaly, which is important to balance the damping
316 effect of atmosphere on the ocean via the surface heat flux (Fig. 3a).

317 Figure 6 provides a schematic summary of the local positive feedback. During
318 the positive phase, the cooling of SST at the North Pacific enhances the meridional
319 gradient of SST south of 40°N. The corresponding increased heat to the south and
320 cooling to the north of the climatological SST front, acts to enhance the gradient of
321 atmospheric temperature and the atmospheric baroclinicity, which favour increased
322 atmospheric transient eddy activity. The increased transient eddy activity and diabatic
323 heating deepen the Aleutian low and drive stronger westerlies. These phenomena
324 result in a large-scale SST cooling in the North Pacific through the ocean dynamical
325 processes. Consequently, this maintains or even enhances the meridional SST gradient
326 anomalies.

327

328 *4.2. Delayed negative feedback in the PDO*

329 The previous section describes the positive feedbacks amplifying PDO
330 anomalies, but a negative feedback is required to transform the PDO to the opposite
331 phase. The PDO pattern shows opposite SST anomalies in the lower latitudes of the
332 Pacific compared to those in the western-to-central North Pacific (Fig. 2). Hence, we
333 focus on the air-sea interaction and upper ocean processes related to the transport of

334 warm water and formation of upper ocean heat content anomalies in the mid-latitude
335 on decadal time scales. All variables are subjected to a 10-yr low-pass Lanczos filter
336 to obtain their decadal components in this section.

337 Figure 7 presents the correlation between the PDO index and the wind stress curl
338 anomalies. At the peak phase of the PDO, the wind stress curl (Fig. 7) shows a dipole
339 pattern in the meridional direction in the central Pacific, with positive values in the
340 subpolar region and negative values in the subtropics. These are consistent with the
341 enhanced surface westerlies during when the PDO is at peak phase (Fig. 5a).
342 Although the significant wind stress curl anomaly appears over the subpolar ocean in
343 ECHAM5/MPI-OM, it does not produce the significantly opposite upper ocean
344 temperature anomaly in the subpolar subsurface, compared to the anomaly in the
345 Kuroshio–Oyashio Extension (KOE) region, which thus does not contribute to a
346 delayed negative feedback. While according to Qiu et al. (2007), the wind stress curl
347 anomaly in the subtropical plays an important role in the phase reversal of the PDO.
348 Thus, a wind stress curl index is defined as the negative of surface wind stress curl
349 anomalies averaged for the area 170°E - 150°W and 20°N - 36°N, which is indicated
350 in Fig. 7a by a green box. To further investigate the evolution of the PDO, we adopt
351 lagged correlation analysis hereafter. The lagged correlation between the PDO index
352 and related variables is described as follows:

$$353 \quad r(m) = \frac{\langle PDO(t)var(t+m) \rangle}{\sqrt{\langle PDO^2(t) \rangle \langle var^2(t+m) \rangle}}$$

354 where *PDO* and *var* indicate the PDO index and related variables, respectively.
355 Angle brackets denote the time average and *m* (>0) denotes the PDO index lead time

356 in years. For example, if $m = 1$ (i.e., the x-coordinate is 1 in Fig. 7b), it means the
357 PDO index leads related variables by one year. Fig. 7b shows that there is a persistent
358 influence of ocean on atmosphere from -1 to +3 years lag, as the correlation between
359 the low-frequency filtered wind stress curl and PDO indices reaches the peak when
360 the wind stress curl index lags the PDO by approximately 1 year.

361 Decadal variability of SST mainly originates from the long-term memory of
362 ocean, particularly from the subsurface process. Thus, analysis of the upper ocean
363 heat content, which includes ocean subsurface signals, is appropriate to understand
364 the mechanisms contributing to decadal variability. To investigate the impact of
365 surface wind stress on the upper ocean, the lagged correlation between the upper
366 700m ocean heat content (OHC700) and the wind stress curl index are displayed in
367 Figure 8. The OHC700 is selected to represent the state of upper ocean, considering
368 the dynamical ocean response to wind forcing as well as the thermocline depth of
369 North Pacific. As Fig. 8a shows, the negative wind stress curl anomalies increase the
370 upper ocean heat content through Ekman convergence in the low latitudes of central
371 Pacific; they exert a continuous influence on the upper ocean caused by persistent
372 anticyclonic wind stress anomalies (Fig. 7). These positive OHC700 anomalies move
373 westward from the ocean interior into the western boundary current (WBC) region
374 (Fig. 8b). When these positive anomalies reach the west Pacific, they warm the WBC
375 region, and are transported by the Kuroshio (Fig 8c). Finally, these positive OHC700
376 anomalies move northward and warm the western-to-central North Pacific (Fig. 8d),
377 and this results in the transition of PDO from positive phase to negative phase after

378 about 18 years. In addition, there is no obvious warm advection of ocean heat content
379 from the subpolar to the KOE region along the Oyashio throughout the entire
380 evolution.

381 To further clarify the role of oceanic Rossby wave in the evolution of the PDO,
382 the sea surface height (SSH) is used to represent the variation of upper ocean. The
383 evolution of SSH anomalies associated with the PDO is shown as the lagged
384 correlation of SSH upon the PDO index (Fig. 9). The positive SSH anomalies appears
385 in the central Pacific at a lag of 0 yr, caused by the anticyclonic wind stress anomalies
386 in the subtropical (Fig. 7). As we move forward from lag 0, the positive SSH anomaly
387 gradually propagates westward, and finally reaches the WBC region at a lag of around
388 5 yr. Subsequently, the negative SSH anomaly appears in the central Pacific.
389 Therefore, considering the phase speed of oceanic Rossby wave in the midlatitude as
390 well as the position of wind stress curl anomaly (Zhang and Delworth 2015, 2016),
391 this phenomenon indicates westward propagation of the OHC700 is caused by free
392 Rossby waves in the subtropical ocean interior.

393 Considering the dominant role of the WBC in the heat transported poleward by
394 the oceans (Kelly and Dong 2004), the lag correlation between PDO index and each
395 term of local heat budget over the upper 700 m is performed to isolate the advection
396 through the Kuroshio in WBC regions (Fig. 10). The local heat budget is described in
397 the following equation:

$$\begin{aligned}
398 \quad \frac{\partial T'}{\partial t} = & - \left(\bar{u} \frac{\partial T'}{\partial x} + u' \frac{\partial \bar{T}}{\partial x} + u' \frac{\partial T'}{\partial x} \right) - \left(\bar{v} \frac{\partial T'}{\partial y} + v' \frac{\partial \bar{T}}{\partial y} + v' \frac{\partial T'}{\partial y} \right) \\
399 \quad & - \left(\bar{w} \frac{\partial T'}{\partial z} + w' \frac{\partial \bar{T}}{\partial z} + w' \frac{\partial T'}{\partial z} \right) + \frac{Q'}{\rho C_p H} + R
\end{aligned}$$

400 where overbars and primes indicate the monthly climatology and anomaly,
401 respectively. The variables u , v , w , and T indicate zonal current, meridional current,
402 vertical velocity, and oceanic temperature averaged over the upper 700 m, respectively.
403 Q is the net surface heat flux into the ocean, ρ is the density of seawater, C_p is the
404 heat capacity of sea water, and R is the residual term. The first four groups of terms on
405 the right-hand side of the equation represent advection in the zonal, meridional,
406 vertical directions, and heat flux, respectively.

407 As Fig. 10a shows, the decadal SST tendency is mainly associated with the
408 meridional temperature advection over the Kuroshio region (indicated by the yellow
409 box in Fig. 8), while heat flux and zonal and vertical temperature advection act as
410 damping terms. A decomposition of the meridional temperature advection term finds
411 that the term related to the mean current and the anomalous temperature gradient
412 $(-\bar{v} \frac{\partial T'}{\partial y})$ is dominant. While over the KOE region (indicated by the green box in Fig.
413 8), the decadal SST tendency is in phase with the zonal temperature advection,
414 whereas heat flux and zonal and vertical temperature advection are out of phase
415 (Fig.10b). A decomposition of the zonal temperature advection term finds that the
416 term related to the mean current and the anomalous temperature gradient $(-\bar{u} \frac{\partial T'}{\partial x})$ is
417 dominant. These results suggest the essential role of climatological current field on
418 the evolution of upper ocean in the WBC regions. The significant warm anomalies

419 appear in the Kuroshio region at a lag of 5 yr, when the downwelling Rossby wave
420 reaches this region, and they reach their peak at a lag of 10 yr. In addition, it is worth
421 noting that there are significant warm anomalies in the KOE region at a lag of 16 yr,
422 and warm anomalies reach the peak at a lag of 18 yr; these are consistent with Figs.
423 8b-c. The above results demonstrate that it takes 18 years to complete the transition
424 from cool phase to warm phase in the KOE region; this is approximately equal to half
425 the PDO period (i.e., about 37 years).

426 To further validate the oscillation on decadal timescales, multichannel singular
427 spectrum analysis (M-SSA) is applied in this study, which is a natural extension of
428 SSA to multiple short noisy time series (Ghil et al. 2002). SSA is mathematically
429 equivalent to the extended EOF analysis, and both are extensions of classical principal
430 component analysis. However, the two methods differ in emphasis: extended EOF
431 analysis usually chooses a number of spatial channels much greater than the number
432 of temporal lags, thus limiting the temporal and spectral information. In MSSA, the
433 situation is the reverse, where the number of temporal lags is large enough to extract
434 detailed temporal and spectral information from the multivariate time series. The
435 advantage over classical spectral analysis is that SSA/M-SSA allows anharmonic
436 behavior to be depicted by a pair of eigenvectors and temporal variations of the
437 amplitude of the oscillatory behavior to be captured.

438 Here, M-SSA is applied to first ten leading PCA components of SST, OHC700
439 and wind stress curl over the North Pacific, and three oscillatory modes can be found
440 on decadal timescales, whose period represents 37, 20 as well as 14 years,

441 respectively. Then, according to Plaut and Vautard (1994), the space-time behavior of
442 the decadal oscillation is investigated by performing a composite analysis keyed to the
443 phases of the oscillation (Fig. 11). In phase 1, negative OHC700 appear over the
444 central and western mid-latitude North Pacific, surrounded by positive anomalies,
445 which indicates the positive phase of PDO. Meanwhile, the anticyclonic wind stress
446 arises over the subtropics and is sustained to phase 2. Through the influence of
447 above-mentioned wind stress, the positive OHC700 arises in the subtropics,
448 propagates westward, and gradually reverses the PDO phase in the following phases,
449 consistent with results of lag correlation. It is worth noting that the cyclonic wind
450 stress is organized at the negative phase of PDO (phase 4). In addition, the evolution
451 of SST is similar that of OHC700, implying the importance of subsurface ocean
452 temperature on the PDO. The OHC700 anomalies over the subpolar gyre region do
453 not seem to be consistent with the wind stress curl, rather these anomalies are
454 consistent with heat flux (Fig. 3a). The wind stress curl anomalies in the subpolar gyre
455 region are also not reversed in phase 4, suggesting that in this region they do not play
456 an important role for the PDO reversal.

457 According to the above-mentioned processes, a decadal time scale, delayed
458 negative feedback for the PDO emerges (Fig. 12). During the positive phase, the
459 cooling of SST at the North Pacific enhances the meridional gradient of SST south of
460 40°N. Therefore, the westerly wind is enhanced and the anticyclonic wind stress
461 anomalies are produced on the south side of westerlies. Consequently, these
462 anticyclonic wind stress anomalies excite the positive OHC700 anomalies by Ekman

463 convergence, and positive OHC700 anomalies propagate westward as the oceanic
464 Rossby wave. When positive OHC700 anomalies reach the western boundary, they
465 warm the WBC region and are transported through the Kuroshio. This eventually
466 leads to warm SST anomalies in the western-to-central North Pacific and reduces the
467 meridional SST gradient south of 40°N, thereby switching the phase of PDO
468 gradually.

469

470 **5. Conclusions and discussion**

471 The PDO is characterized by a horseshoe pattern of SST anomalies with one sign
472 in the western-to-central North Pacific surrounded by anomalies of opposite sign. It is
473 an important climate variation on decadal timescales that has a major influence on
474 climate system. Understanding the mechanism of the PDO is therefore of broad
475 scientific and socioeconomic interest, but understanding remains limited, especially of
476 the role of extra-tropical ocean–atmosphere interaction in the PDO. This paper
477 reinvestigates the mechanisms for the PDO focusing on this last aspect and using a
478 150-year preindustrial control simulation of ECHAM5/MPI-OM, which reproduces
479 key features of the PDO very well. The simulated positive (negative) PDO exhibits a
480 large-scale SST cooling (warming) over the central and western mid-latitude North
481 Pacific, and the atmospheric geopotential height shows an equivalent barotropic low
482 (high) anomaly throughout the troposphere associated with the enhanced (weakened)
483 zonal wind. This equivalent barotropic cold/trough structure is the typical structure of
484 the observed PDO during winter. Based on the meridional SST gradient in the KOE

485 region, we identify two possible feedbacks – a positive and a delayed negative – that
486 may sustain the PDO.

487 The positive feedback process to maintain the structure of PDO is as follows.

488 During the positive (negative) PDO phase, the cooling (warming) of SST at the North

489 Pacific enhances (reduces) the meridional SST gradient south of 40°N. The

490 subsequent change in net surface heat flux damps the meridional gradient of SST, and

491 provides the diabatic heat driving the increased (decreased) meridional gradient of

492 atmosphere temperature as well as baroclinicity in mid-to-lower levels. This favours

493 (suppresses) transient atmospheric variability and diabatic heating. The atmospheric

494 response to either shallow diabatic heating (cooling) or deep transient eddy heating

495 (cooling) is dynamically baroclinic, whereas the response to the transient eddy

496 vorticity forcing is equivalent barotropic. These intensify (weaken) the Aleutian low.

497 The stronger westerlies result in the ocean dynamical processes, which is important to

498 balance the damping effect of surface heat flux on the ocean. Consequently, they will

499 reinforce the enhanced meridional SST gradient south of 40°N.

500 The delayed negative feedback reversing the PDO phase on decadal timescales is

501 summarized as follows. During the positive (negative) phase, the SST is cold (warm)

502 in the North Pacific and the meridional SST gradient is intensified (weakened) south

503 of 40°N. Therefore, the westerly wind is intensified (reduced) and the anticyclonic

504 (cyclonic) wind stress anomalies are produced on the south side of westerlies. Then,

505 these anticyclonic (cyclonic) wind stress anomalies excite the warm (cold) upper

506 ocean anomalies by Ekman convergence (divergence). These upper ocean anomalies

507 propagate westward from the ocean interior into the WBC region, with a speed
508 consistent of free Rossby waves. When they reach the western boundary, they are
509 transported by the Kuroshio as well as the KOE, and reverse the sign of the
510 preexisting PDO and the meridional SST gradient gradually.

511 Our conclusion that mid-latitude ocean-atmosphere interaction can play an
512 important role in the evolution of PDO is consistent with other studies (e.g., Latif
513 2006; Latif and Barnett 1994, 1996; Park and Latif 2010; Zhang and Delworth 2015).
514 However in contrast to these, our work clearly shows that the meridional SST gradient
515 in the KOE region plays an essential role in the positive and the delayed negative
516 feedbacks, due to its influence on the internal atmospheric variability through not only
517 direct diabatic heating but indirect transient eddy forcing. Although the feedback
518 proposed in this study is based on the results from ECHAM5/MPI-OM, the types of
519 feedback coincide with those in observations (Fang and Yang 2016; Latif and Barnett
520 1996; Qiu et al. 2007). The positive and negative feedback constitute a complete
521 evolution of the atmosphere and ocean that is useful for understanding the oscillatory
522 mechanism of the PDO. The results here are obtained from only one single model,
523 and the robustness of feedbacks identified here should be confirmed in other models.
524 In particular, it is important to understand whether the strength of the feedbacks is
525 realistically reproduced. This question is difficult to address and beyond the scope of
526 this work. However, there is some evidence from models with higher resolution than
527 used here (e.g., Ma et al. 2015; Ma et al. 2016; O'Reilly and Czaja 2015; Zhou et al.
528 2015) and climate predictions (Scaife and Smith 2018) that the strength of

529 mid-latitude ocean-atmosphere interaction in current climate models is
530 underestimated.

531 In addition, some other important mechanisms on modulating PDO, such as the
532 atmospheric bridge (Alexander et al. 2002; Trenberth et al. 1998) and the
533 extratropical–tropical interaction (Gu and Philander 1997; McPhaden and Zhang
534 2002), are not discussed in this paper. In the present study, we only focused on the
535 influence of mid-latitude air-sea interaction on the PDO. Therefore, assessing the
536 relative importance of various mechanisms would be an important aspect in future
537 works. Besides, climate variability is influenced by initial conditions as well as
538 external forcing (e.g., Meehl et al. 2014), and data assimilation can provide better
539 constraints on initial conditions for climate simulations and predictions (e.g.,
540 Keenlyside et al. 2005; Keenlyside et al. 2008; Zheng and Zhu 2010; Zheng et al.
541 2007; Zheng et al. 2006). Thus, in a following work we will conduct a quantitative
542 study via the established data assimilation system (Luo et al. 2017) to confirm the
543 relevance of the mechanism of this study. For example, we can reproduce the
544 meridional SST gradient in the KOE region, by assimilating SST, to investigate to
545 what extent the atmospheric anomalies are reproduced under this SST gradient
546 structure. And we can reconstruct the variability of tropical Pacific, through only
547 assimilating SST in the tropical Pacific, to quantify the contribution of tropical Pacific
548 on the PDO and to explore underlying mechanisms.

549

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557 data are available from
558 https://www.esrl.noaa.gov/psd/data/gridded/data.20thC_ReanV2c.html.

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820

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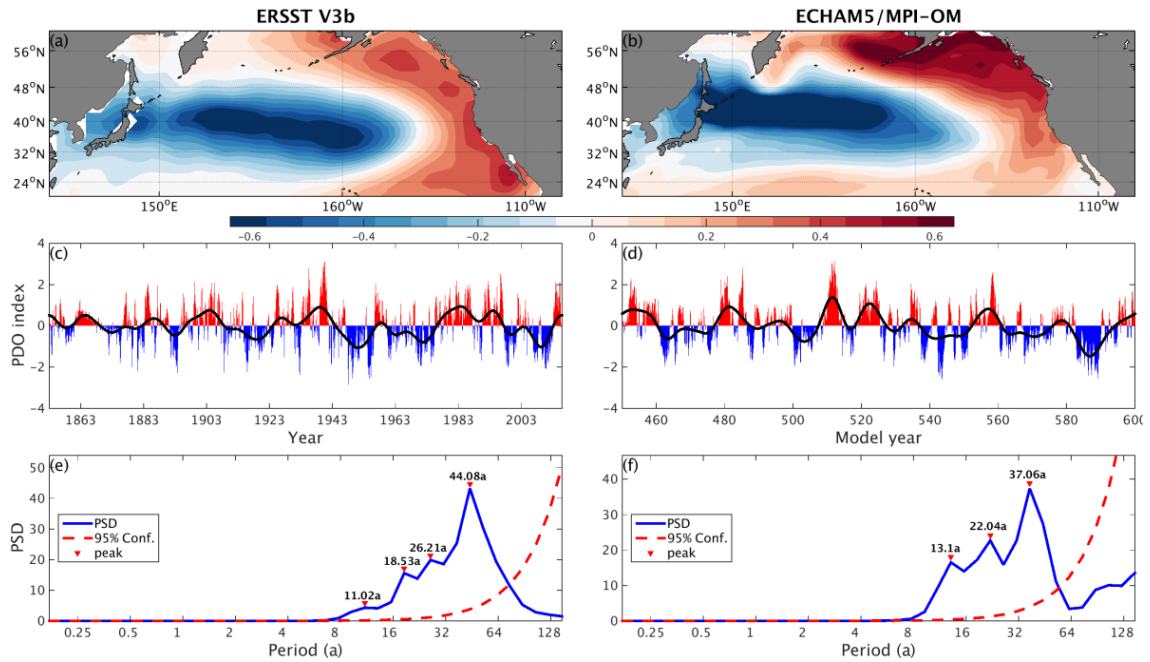
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Figure 1



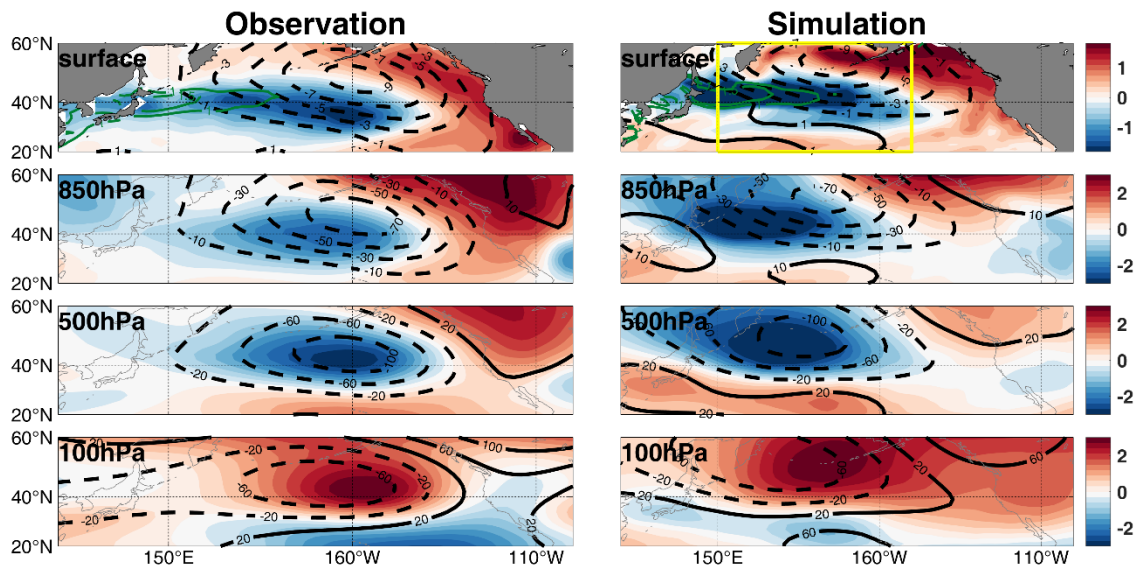
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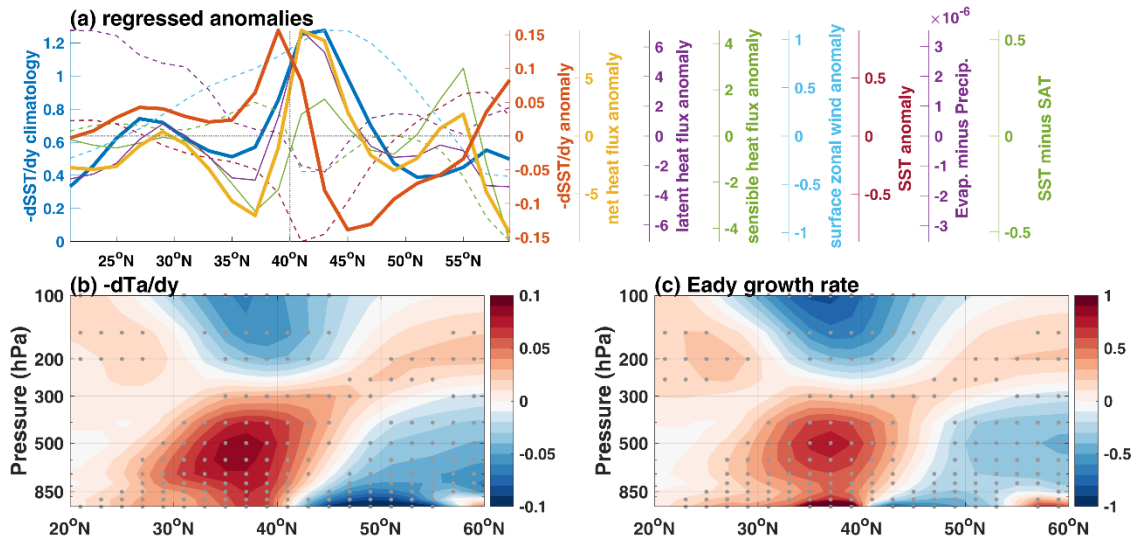
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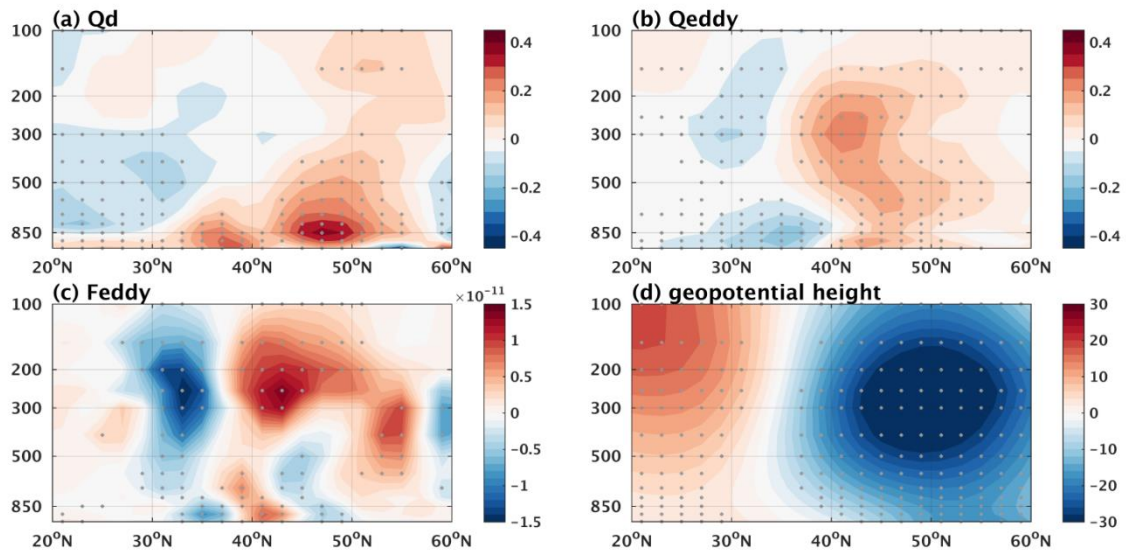
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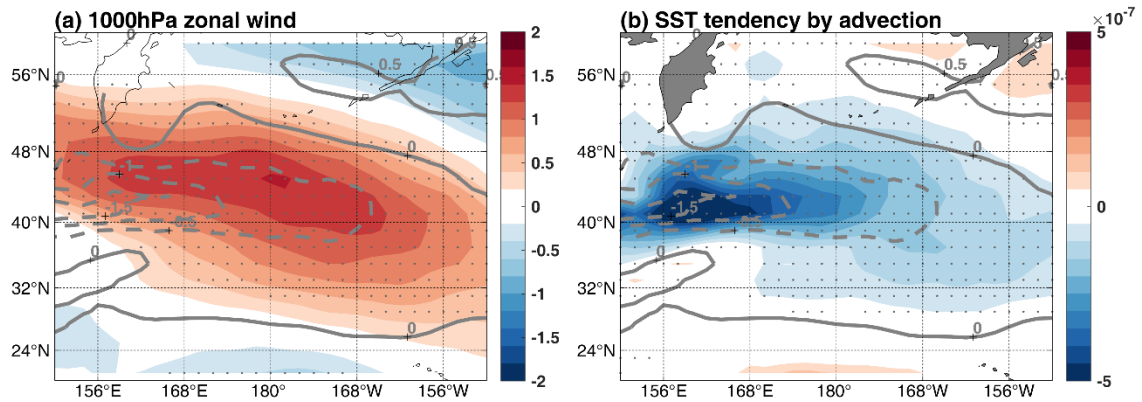
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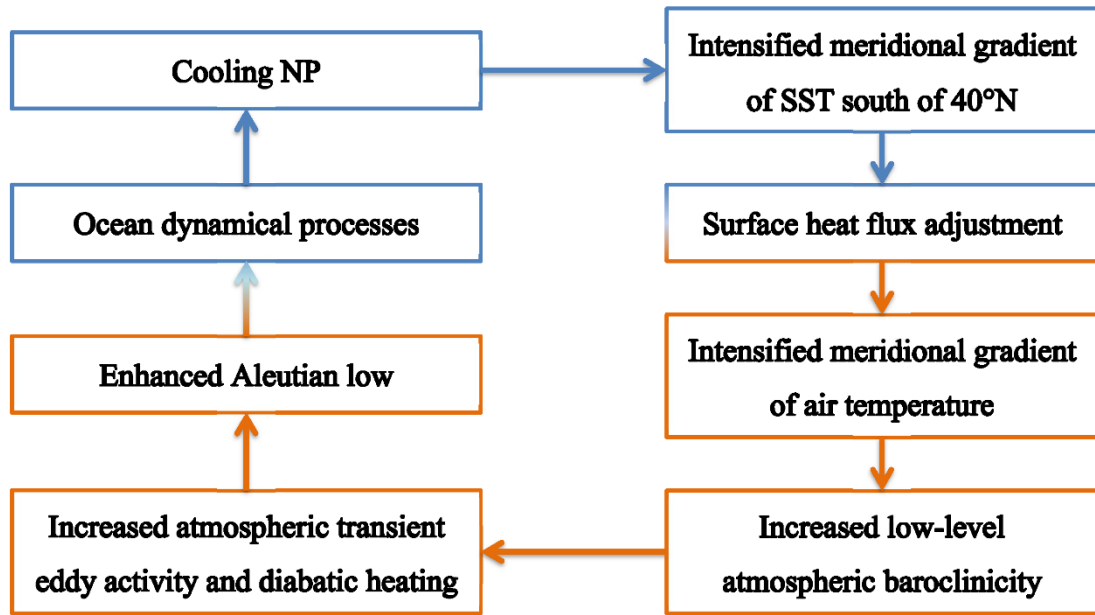
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Figure 6



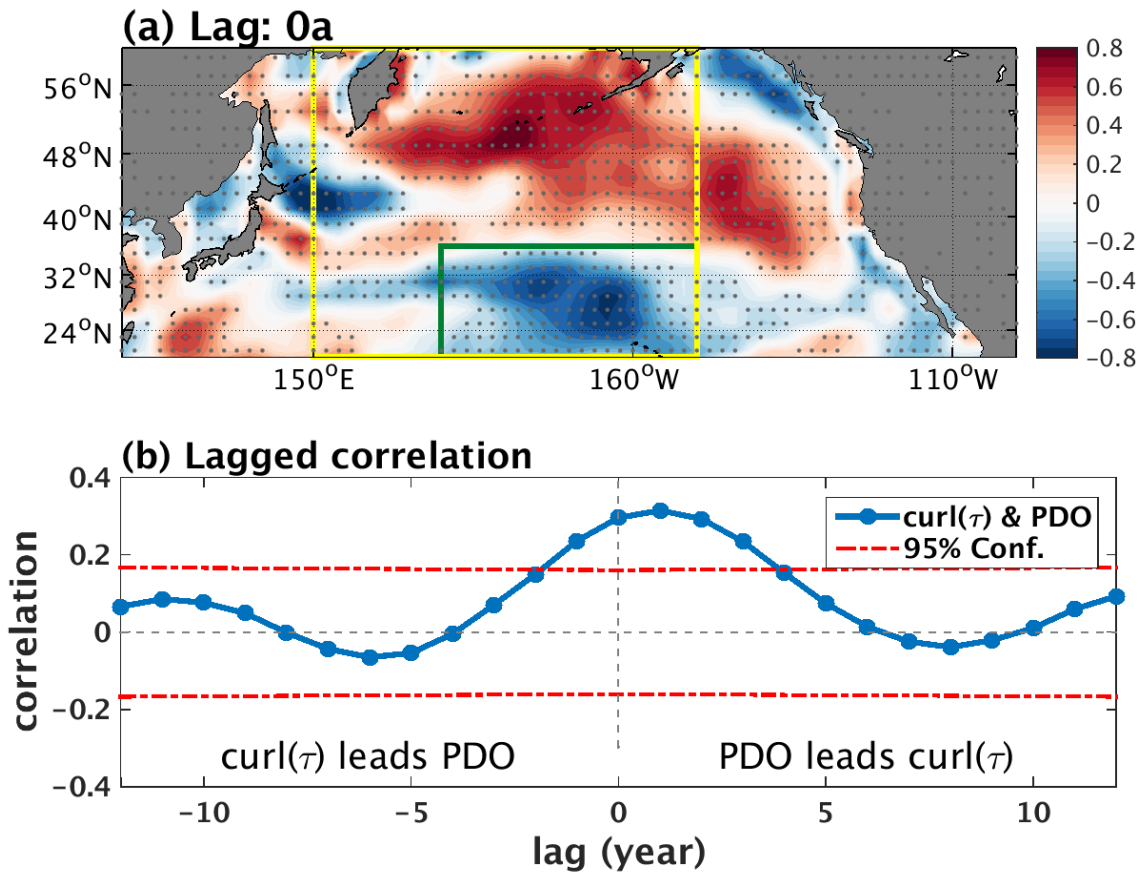
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Figure 7



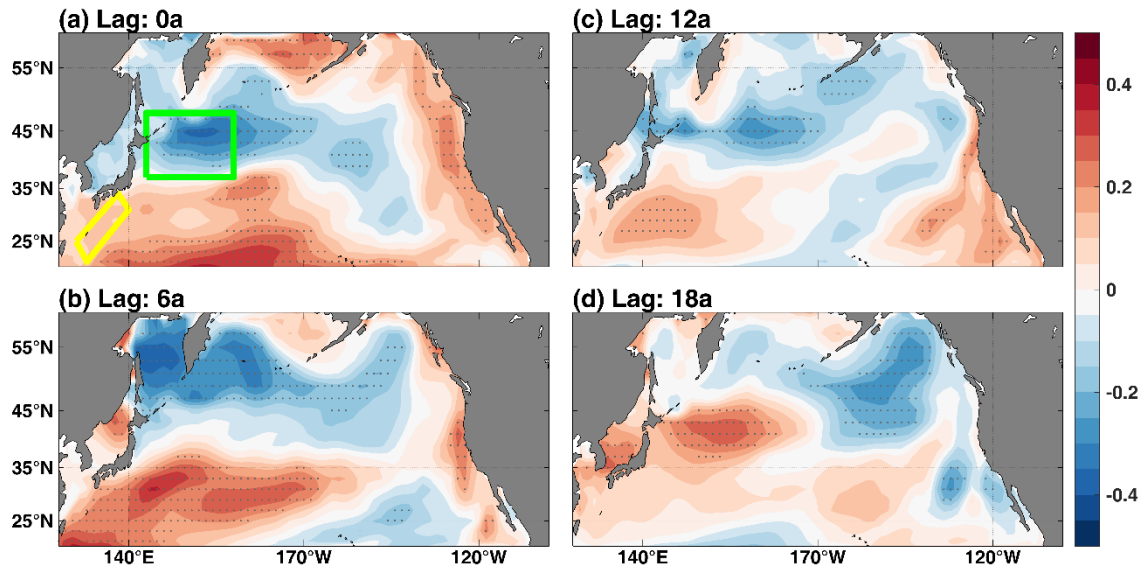
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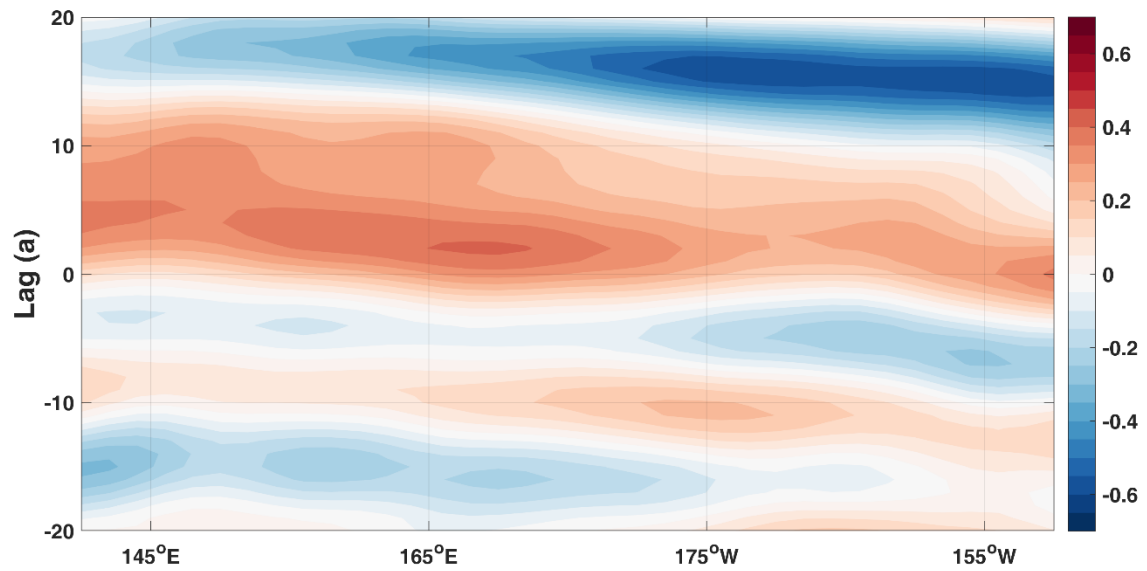
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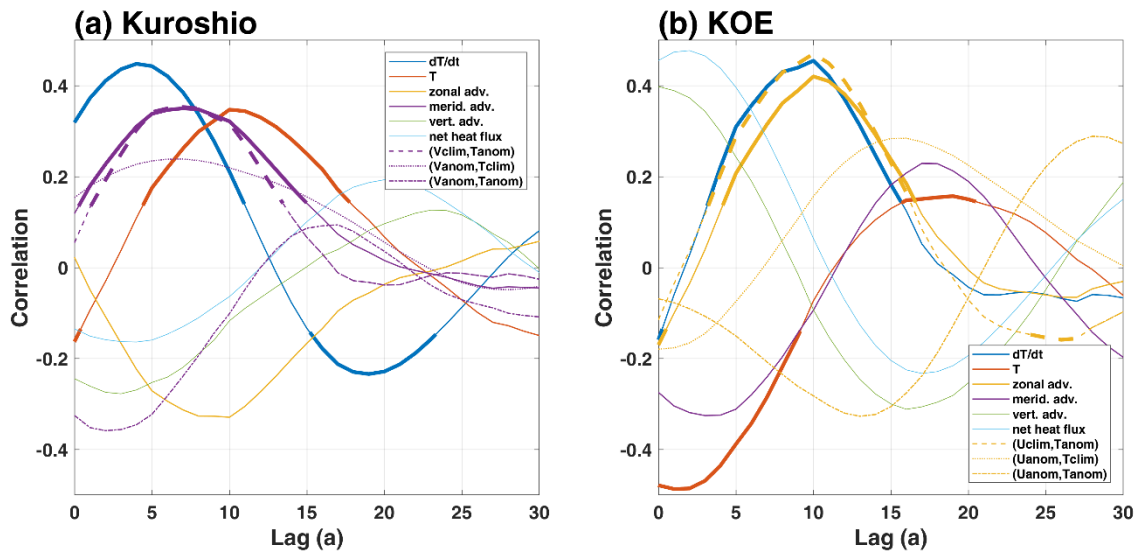


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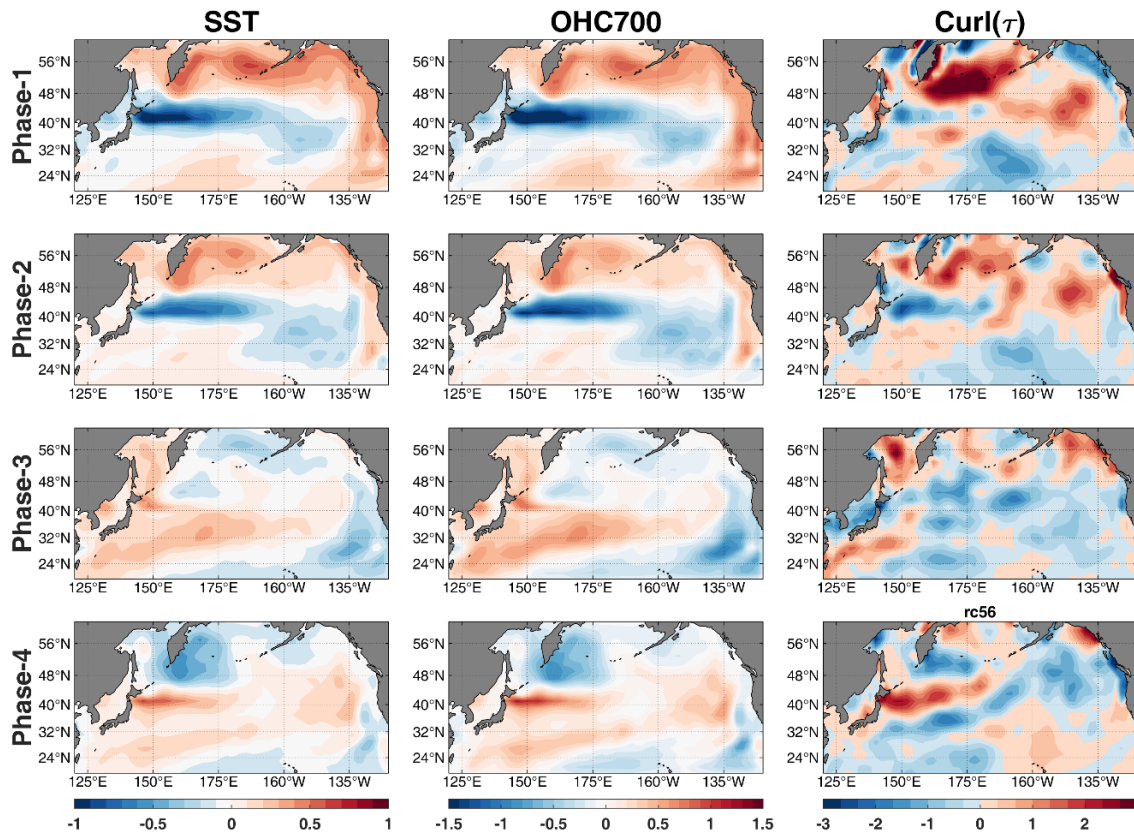


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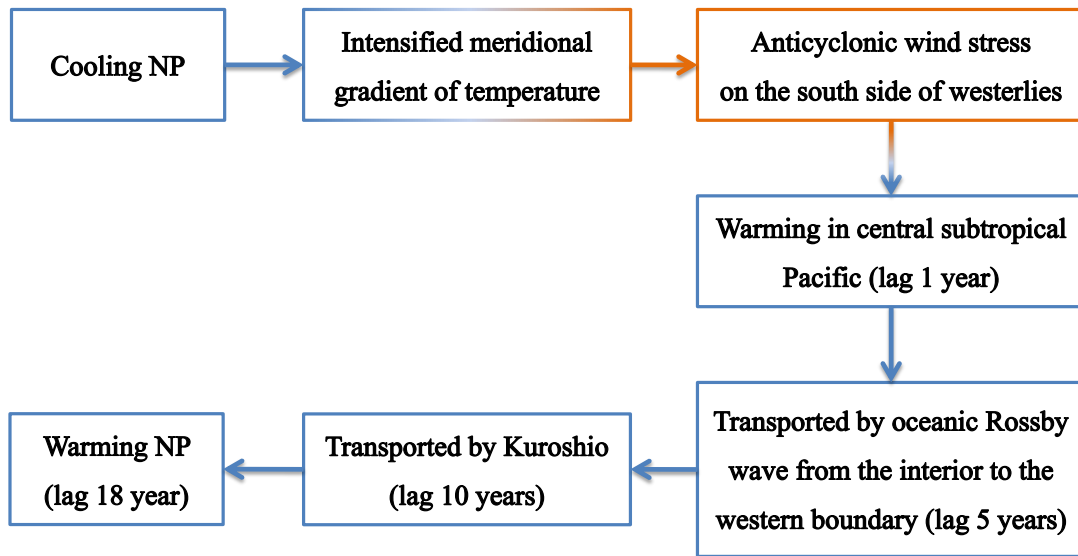
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