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# Microfluidic hydrogen storage capacity and residual trapping during cyclic injections: Implications for underground storage

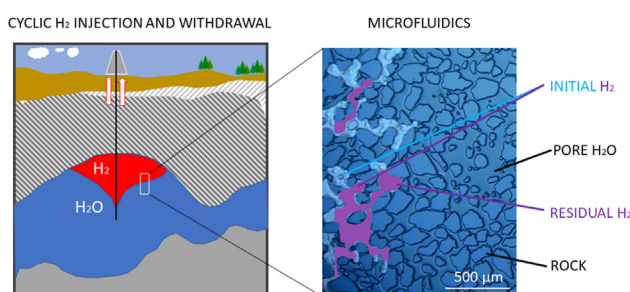
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## HIGHLIGHTS

- Microfluidics for examination of cyclic hydrogen injections.
- Microscopic hydrogen storage capacities up to 60% of the pore space.
- Reproducible residual hydrogen saturation between the injection cycles.
- Efficient reconnection of residually trapped hydrogen in the next drainage cycle.

## GRAPHICAL ABSTRACT



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## ABSTRACT

Long-term and large-scale  $H_2$  storage is vital for a sustainable  $H_2$  economy. Research in underground  $H_2$  storage (UHS) in porous media is emerging, but the understanding of  $H_2$  reconnection and recovery mechanisms under cyclic loading is not yet adequate. This paper reports a qualitative and quantitative investigation of  $H_2$  reconnection and recovery mechanisms in repeated injection-withdrawal cycles. Here we use microfluidics to experimentally investigate up to 5 cycles of  $H_2$  injection and withdrawal under a range of injection rates at shallow reservoir storage conditions. We find that  $H_2$  storage capacities increase with increasing injection rate and range between ~10% and 60%. The residual  $H_2$  saturation is in the same range between cycles (30–40%), but its distribution in the pore space visually appears to be hysteretic. In most cases, the residually trapped  $H_2$  reconnects in the subsequent injection cycle, predominantly in proximity to the large pore clusters. Our results provide valuable experimental data to advance the understanding of multiple  $H_2$  injection cycles in UHS schemes.

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## Introduction

Hydrogen ( $H_2$ ) is an emission-free energy carrier and its wider use can contribute to climate change mitigation by decreasing the share of fossil fuels in the global energy mix. Full-scale industrial implementation in a global  $H_2$  economy will require numerous storage sites and solutions [1]. Future  $H_2$  storage demand in Europe is predicted to range between 63 and 180 billion standard  $m^3$  in 2050, assuming  $H_2$  total demand of 780–2251 TWh [2] and 24% storage capacity [3]. Underground  $H_2$  storage (UHS) in depleted hydrocarbon fields and aquifers has been proposed as a reliable and safe storage technology due to the presence of an impermeable seal and large pore space [4,5]. The technicalities are similar to natural gas storage (UGS), where cushion gas remains in reservoir to maintain the target pressure and working gas is injected at peak supply (summer) and withdrawn at peak demand (winter). However,  $H_2$  is a low density and low viscosity gas with high diffusivity and biogeochemical reactivity and therefore its behavior in porous media will differ from that of other gases. The experience with UHS in porous media is limited to: 1) two pilot tests in depleted gas fields [6,7] and 2) town gas storage in aquifers [8,9].

Scientific challenges relevant for the UHS arise from  $H_2$  physical properties as well as reactions with rock minerals and microorganisms, potentially reducing the storage efficiency [5,10].  $H_2$  injections are prone to unstable displacement and gravity override due to low viscosity and density. Moreover, a certain amount of  $H_2$  may be permanently lost during storage operation by various physical, chemical and operating loss mechanisms [4]. Reservoir simulation and wettability studies are currently dominating the research literature in the field of UHS flow physics. Storage capacities and recovery factors have been estimated using conceptual reservoir models with extrapolated input parameters not specifically measured for  $H_2$  [11,12].

Contact angle measurements indicated that  $H_2$  is a non-wetting fluid on pure quartz surfaces and sandstones but shifts to intermediate-wet state in the presence of organic acids [13–16]. No clear difference in contact angles was reported between  $H_2$ ,  $CH_4$ ,  $H_2$ – $CH_4$  mixtures and  $N_2$  using a captive-bubble method in sandstones and a borosilicate micromodel at pressures between 10 and 100 bar [17,18]. On the other hand,  $CO_2$  was found to be more wetting, i.e. higher contact angles, compared to  $H_2$  [13,16]. Note that the contact angle differences between various gases are expected to become more pronounced at pressures above 100 bar due to the increasing gas density differences, shown for basaltic rocks [19]. Moreover, rock-gas interfacial tension calculations indicated significant differences between  $H_2$ ,  $CH_4$  and  $CO_2$  [20]. There are also discrepancies regarding the influence of pressure, temperature and salinity on  $H_2$  wettability.  $H_2$  contact angles increased with increasing pressure and temperature when using the tilted plate method [13,16], whereas no meaningful effect of pressure, temperature and salinity was reported using the captive bubble technique under a range of 7–207 bar, 20–50 °C and 1000–50000 ppm NaCl brine [14,15,17].

Neither reservoir simulations nor wettability studies can adequately describe pore scale influencing factors on UHS

such as interactions between  $H_2$ , reservoir rock and its native fluid as well as trapping mechanisms [21]. Residual trapping is recognized as one of the major  $H_2$  loss mechanisms [4,22], which is expected to decrease with decreasing capillary forces and increasing  $H_2$  wetting [23]. The UHS involves multiple cycles of  $H_2$  injection (drainage) and withdrawal (imbibition), and residual trapping occurs during imbibition where water is available and mobile in the reservoir, e.g. in the  $H_2$ – $H_2O$  transition zones. The residually trapped  $H_2$  ganglia may reconnect during drainage due to hysteresis. Hysteresis was evident from relative permeability measurements [24,25] and microfluidic-based contact angle measurements [26], but it remains unaddressed for several injection-withdrawal cycles.

A few laboratory investigations of  $H_2$  residual trapping used in-situ visualization of a single drainage-imbibition cycle in different sandstones. The initial and residual  $H_2$  saturations in the pore space (values between 0 and 1) were measured to be 0.65 and 0.41, respectively [27].  $H_2$  recovery decreased from 43.1% when flooded with non- $H_2$ -equilibrated brine to 31.6% for  $H_2$ -equilibrated brine [28]. The initial  $H_2$  saturation was ~6 times lower compared to  $N_2$  using the same injection rate [29]. No clear pressure impact on the initial  $H_2$  saturation was evident, contrary to  $H_2$  residual trapping which increased with increasing pressure and decreasing injection rate [30].

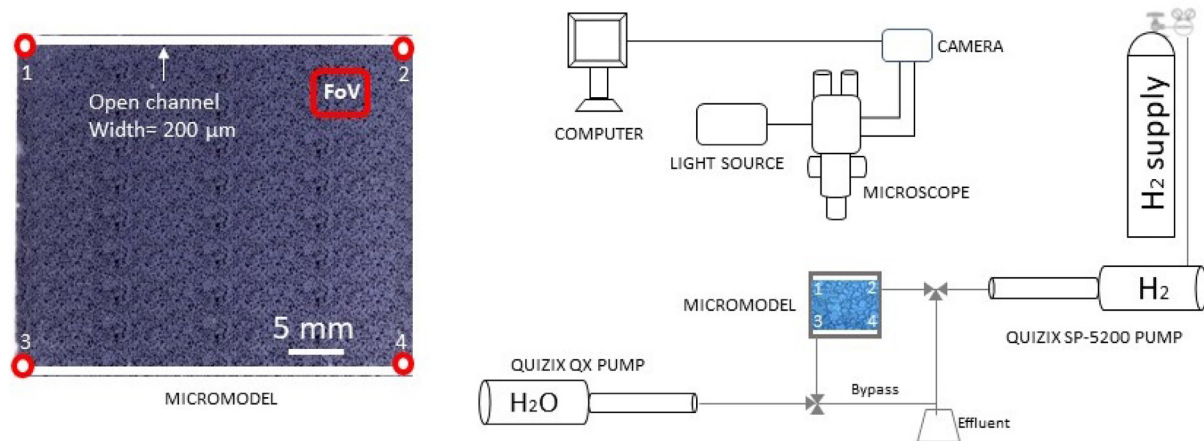
$H_2$  cyclic injections were only performed for two drainage-imbibition cycles and resulted in similar initial and residual  $H_2$  saturations: 0.48 and 0.07, respectively [30]. In contrast, cyclic injections have been extensively investigated for  $CO_2$  storage, where some studies indicated an increased residual trapping over the injection cycles [31–33], contradicting a classic trapping theory [34]. More systematic studies with the increased number of injection cycles are required to find out whether the multiple injection cycles can potentially result in hysteresis and increase  $H_2$  residual trapping.

Microfluidics is a valuable tool for direct observations of pore space to corroborate core scale measurements. Small pore volumes are beneficial for the sake of time and safety when working with highly flammable  $H_2$  gas at elevated pressures. In this work, we qualitatively describe hysteretic  $H_2$  trapping and reconnection mechanisms during cyclic injections in a silicon-wafer micromodel with pore patterns resembling a natural sandstone. Up to 4–5 drainage-imbibition cycles were run under a wide range of injection rates at 40 bar and ambient temperature, representing the storage conditions of a shallow aquifer or a gas-water transition zone in a depleted gas field. An in-house MATLAB code was developed to quantify microscopic  $H_2$  storage capacity, residual trapping and recovery factors. Our results add new experimental data, enhancing the understating of hysteretic  $H_2$  behavior during multiple injection cycles.

## Materials and methods

### Porous material

We used a silicon micromodel capable of withstanding pressure up to 150 bar (Fig. 1). The irregular sandstone-based pore patterns (pure quartz) were etched and repeated 36 ( $4 \times 9$ ) times on the silicon wafer using deep reactive ioning etching



**Fig. 1** – Left: Micromodel with irregular pore patterns replicating natural sandstone. The micromodel was built with four ports (1–4) and two open channels (from the ports 1 to 2 and from 3 to 4). The field of view (FoV) refers to the micromodel area observed by the microscope (not to scale). Right: Experimental setup consisting of the micromodel, two pumps for H<sub>2</sub> and H<sub>2</sub>O injections and pressure control as well as the microscope equipped with the light source and the camera. H<sub>2</sub> and H<sub>2</sub>O were injected diagonally from the opposite ports, i.e. from ports 2 and 3, respectively.

(DRIE) with an etching depth of 30  $\mu\text{m}$ . The DRIE technique ensured a correct reproduction of morphological and topological features, preserving the sharp grain walls with a surface roughness of 100 nm and high aspect ratio and coordination number, which ensured a correct magnitude of the capillary forces. The micromodel top (transparent borosilicate glass) and bottom (silicon) surfaces were anodically bonded and produced with strongly hydrophilic surfaces, with measured H<sub>2</sub> contact angles ranging between 19° and 60°. The micromodel surfaces were not aged in organic acids, making them more hydrophilic than natural reservoirs [13,16]. The micromodel has two open channels (100% void space), connected from ports 1 to 2 and from 3 to 4. The pore network has a length of 27 mm and a width of 21.4 mm, with a total porosity of 61% (Table 1). The micromodel studied region is defined as the field of view (FoV) and represents approximately 1% of the entire area micromodel. A more detailed description of the micromodel construction procedure and its properties can be found elsewhere [35,36].

### Experimental setup and procedures

The micromodel was assembled in the PEEK holder with four outlet ports connected to the 1/16" PEEK tubing, where two tubes from the diagonally located ports (2 and 3) were connected to two Quizix pumps through 1/16" stainless-steel tubing (Fig. 1). Quizix QX pump was filled with filtered

deionized H<sub>2</sub>O, whereas H<sub>2</sub> was accommodated by Quizix SP-5200 pump (cylinder C5000-10K-SS-AT). A microscope (Nikon SMZ1500) connected to a camera (Nikon D7100) and computer enabled us to directly observe the micromodel FoV, which was illuminated by a light source with spot lighting (Photonic LED F1 Cold light 5500K).

Two different groups of experiments were performed at pore pressure of 40 bar and ambient temperature ( $20 \pm 1$  °C): 1) Single-cycle of H<sub>2</sub> injection and withdrawal, i.e. primary drainage and imbibition only (experiments A1–A4) and 2) multiple cycles of H<sub>2</sub> injection and withdrawal (experiments B1–B4). An overview of experiments and key results are shown in Table 3. In the experiments A1–A4, the H<sub>2</sub>O pump was used for H<sub>2</sub>O withdrawal (drainage) and injection (imbibition), whereas the constant pressure in the micromodel was maintained by the H<sub>2</sub> pump. In total, single-cycle injection-withdrawal experiments were performed four times at different injection rates in the range of 0.1–50 mL/h.

In experiments B1–B4, the pump operation modes were different. During drainage, the H<sub>2</sub>O pump was set to constant pressure and H<sub>2</sub> was injected from the H<sub>2</sub> pump at constant flow rate. After drainage, the tubing connection from the H<sub>2</sub>O pump to the micromodel was cleaned with H<sub>2</sub>O via the by-pass tubing to remove the remaining H<sub>2</sub>, preventing the H<sub>2</sub>–H<sub>2</sub>O slug flow in the micromodel. During imbibition, both pumps were operated at constant flow rates where H<sub>2</sub>O was injected in the micromodel while the piston in the H<sub>2</sub> pump retracted. Imbibition was terminated after the establishment of the residually trapped H<sub>2</sub> ganglia. Then the system was ready for a new drainage-imbibition cycle, which was repeated three-four times. Prior to a new cycle, H<sub>2</sub> was injected to the bypass line to remove the remaining H<sub>2</sub>O. In total, four cyclic experiments were run at various injection rates in the range of 1–10 mL/h.

Dimensionless numbers describe the interplay between various forces acting on two-phase flow. In this work, the capillary number (the ratio of viscous to capillary forces) is defined as  $N_{Ca} = U \cdot \mu / \sigma$ , where  $U$  is the injection velocity [m/s],  $\mu$  is the invading fluid viscosity [ $\mu_{H_2} = 8.8 \times 10^{-6}$  Pa s and

**Table 1** – Micromodel properties.

	Micromodel	Micromodel FoV
Length [mm]	27	3.5
Width [mm]	21.4	1.96
Depth [ $\mu\text{m}$ ]	30	
Pore volume [ $\mu\text{L}$ ]	11	0.09
Porosity [frac.]	0.61	0.44
Permeability [D]	2.97	
Pore throat length [ $\mu\text{m}$ ]	10–300	

**Table 2 – Flow conditions: Injection rate (Q) and injection velocity (U); and dimensionless numbers: Capillary ( $N_{Ca}$ ), Reynolds (Re), Peclet (Pe), and Bond numbers (Bo).**

Q [mL/h]	U [m/day]	$N_{Ca}$		Re		Pe	Bo
		Drainage	Imbibition	Drainage	Imbibition		
0.1	4.9	$6.8 \times 10^{-9}$	$7.7 \times 10^{-7}$	0.002	0.006	1.6	0.0017
1	48.6	$6.8 \times 10^{-8}$	$7.7 \times 10^{-6}$	0.02	0.06	15.9	
2.5	121.4	$1.7 \times 10^{-7}$	$1.9 \times 10^{-5}$	0.06	0.16	39.7	
5	242.9	$3.4 \times 10^{-7}$	$3.8 \times 10^{-5}$	0.12	0.32	79.3	
10	485.7	$6.8 \times 10^{-7}$	$7.7 \times 10^{-5}$	0.23	0.63	158.6	
50	2428.7	$3.4 \times 10^{-6}$	$3.4 \times 10^{-4}$	1.16	3.17	793.2	

**Table 3 – Initial ( $S_{gi}$ ) and residual ( $S_{gr}$ )  $H_2$  saturations and recovery factors, defined as  $(S_{gi} - S_{gr})/S_{gi}$ , during single-cycle (A1–A4) and multiple-cycles (B1–B4) injections.**

Exp ID	Q [mL/h]	Cycle number	$S_{gi}$ [fraction]	$S_{gr}$ [fraction]	Recovery factor [fraction]
A1	0.1	1	0.09	0.04	0.53
A2	1	1	0.18	0.05	0.71
A3	10	1	0.61	0.33	0.45
A4	50	1	0.47	0.30	0.38
B1	1	1	0.14	0.03	0.44
		2	0.09	0.04	0.49
		3	0.13	0.07	0.48
		4	0.06	0.08	N/A
B2	2.5	1	0.36	0.35	0.02
		2	0.45	0.40	0.10
		3	0.50	0.42	0.16
		4	0.38	0.33	0.17
B3	5	5	0.73	0.48	0.34
		1	0.42	0.32	0.23
		2	0.60	0.31	0.48
		3	0.67	0.28	0.59
B4	10	4	0.53	0.32	0.40
		5	0.60	0.34	0.42
		1	0.50	0.29	0.42
		2	0.56	0.31	0.45
		3	0.42	0.29	0.32
		4	0.62	0.27	0.57
		5	0.58	0.28	0.51

$\mu_{H_2O} = 1.0 \times 10^{-3}$  Pa s [37]), and  $\sigma$  is the  $H_2$ – $H_2O$  interfacial tension [ $= 0.073$  N/m [38]]. The injection velocity was calculated as follows:  $U = Q/(L \cdot d \cdot \phi)$ , where Q is the injection rate [ $m^3/s$ ],  $\phi$  is the micromodel porosity [fraction], and L and d are the micromodel length and depth [m], respectively. The Reynolds number (the ratio of inertial to viscous forces) is defined as  $Re = \rho \cdot U \cdot D_{50} / \mu$ , where  $\rho$  is the invading fluid density [ $\rho_{H_2} = 3.2$  kg/ $m^3$  and  $\rho_{H_2O} = 1000$  kg/ $m^3$ ] and  $D_{50}$  = median grain diameter [ $= 1.1 \times 10^{-4}$  m] – an approximation of the characteristic length scale [39]. The Peclet number correlates convection and diffusion transport and is defined as  $Pe = U \cdot D_{50} / D$ , where D is the  $H_2$  diffusion coefficient through water equal to  $4 \times 10^{-9}$   $m^2/s$  [40]. The Bond number (the ratio of gravitational to surface tension forces) is defined as  $Bo = \Delta \rho \cdot g \cdot (D_{50})^2 / \sigma$ , where  $\Delta \rho$  is the density difference between  $H_2$  and  $H_2O$ , and g is the acceleration due to gravity. The range of various dimensionless numbers (Table 2) was estimated at experimental conditions and indicated that the  $H_2$ – $H_2O$  flow occurred under the laminar flow regime, with the dominance

of convection and surface tension (i.e., capillary forces) over diffusion and gravity. The interplay between viscous and capillary forces was non-trivial where both forces could compete because the experimental  $N_{Ca}$ -range belongs to the transition zone in the Log ( $N_{Ca}$ )-flow diagram [26,41].

### Image analysis

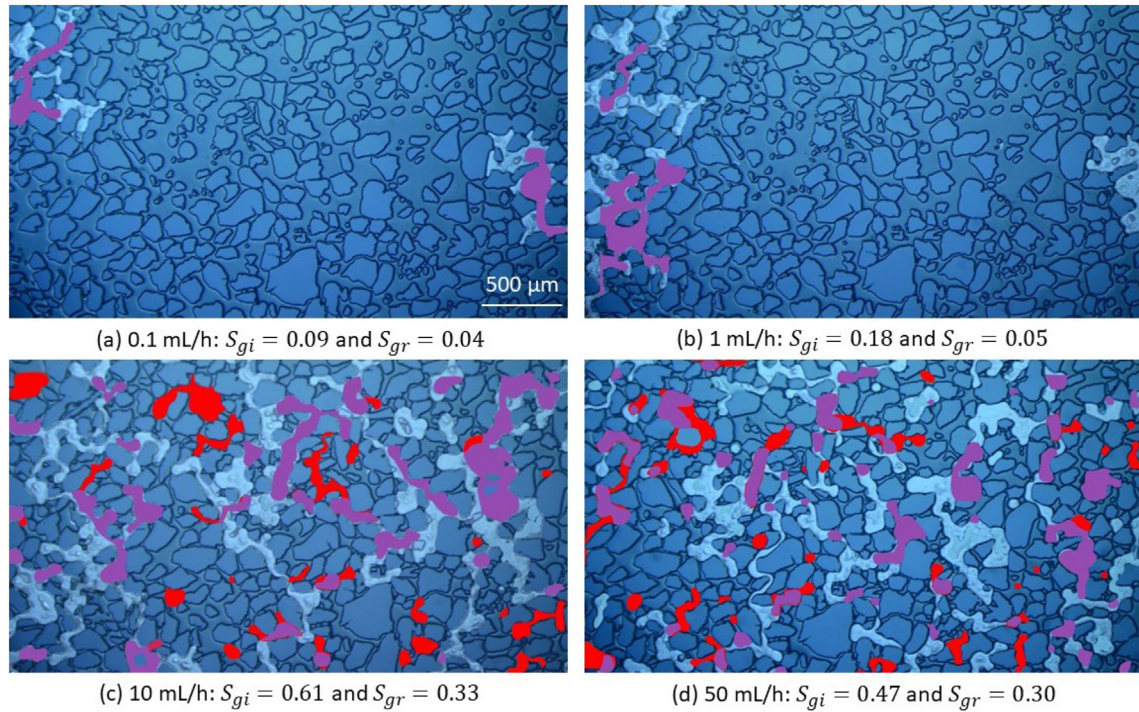
The raw images were processed and analyzed to calculate the FoV porosity and  $H_2$  saturation using a combination of an open-source ImageJ software and in-house MATLAB code. The color gradients due to a spotlight required the image pre-processing with manual segmentation of the grains. The FoV porosity was therefore calculated for each image using color thresholding in ImageJ before further analysis in MATLAB. The  $H_2$  saturations were calculated based on the in-house MATLAB code that used the background subtraction algorithm, with a background image of 100%  $H_2O$ -saturated FoV. The average relative uncertainty of  $H_2$  saturation was estimated to be 9% and was related to the noise threshold, caused by inclusion of the  $H_2O$  droplets and small grains in the  $H_2$  saturation. By adjusting a threshold value of several sequential images with equal quasi-steady-state  $H_2$  saturation, the relative uncertainty was calculated as standard deviation.

## Results and discussion

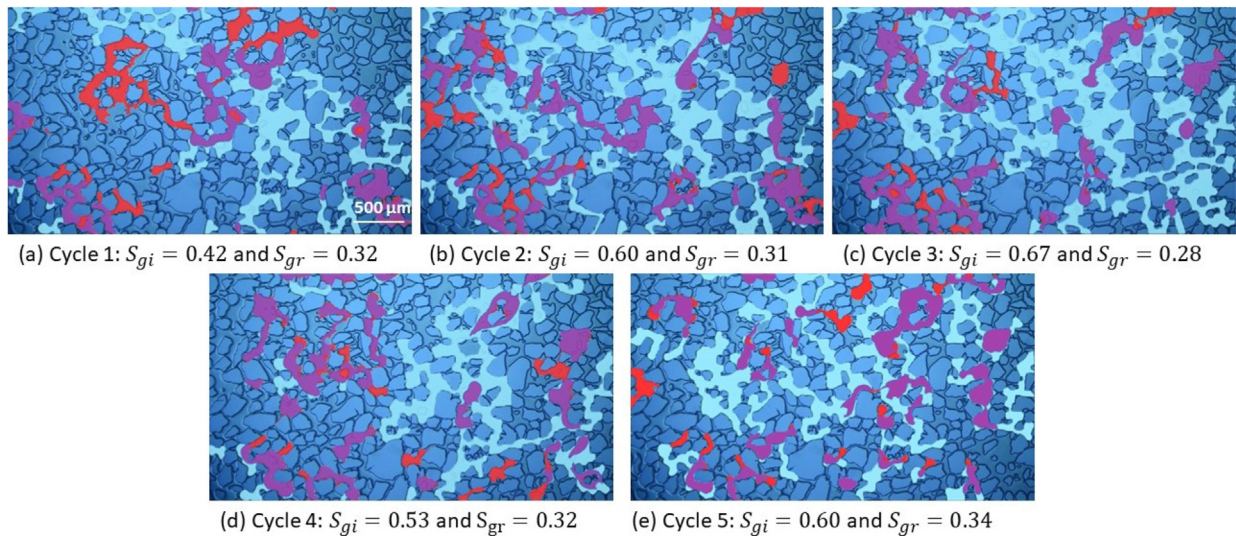
### Displacement, trapping and re-connection mechanisms

Primary drainage injections at low rates ( $\leq 1$  mL/h) resulted in the low  $H_2$  saturation in the FoV ( $S_g < 0.20$ ) due to high capillary entry pressures (Fig. 2, Table 3). At high injection rates ( $\geq 10$  mL/h) the  $H_2$  saturation increased by  $\sim 2$ – $3$  times and both connected and disconnected  $H_2$  established due to Roof snap-off [42].  $H_2$  displacement and trapping during imbibition was governed by I1 imbibition and I2 imbibition mechanisms, respectively [43].  $H_2$  was displaced from several pores to a single pore (I1 imbibition), where  $H_2$  was disconnected at the pore wall and residually trapped (I2 imbibition). Distribution of the residually trapped  $H_2$  after imbibition (red + purple in Fig. 2) depended on the initial  $H_2$  distribution after drainage (blue + purple in Fig. 2). In most cases, the residually trapped  $H_2$  remained in the same pores (purple in Fig. 2c and d), but displacement to the neighboring pores was also observed (red in Fig. 2c and d). The observed displacement and trapping mechanisms corroborated our previous study in the same micromodel at 5 bar [26], suggesting that displacement





**Fig. 2** – Combined images of  $H_2$  saturation after primary drainage,  $S_{gi}$  (blue + purple), and after imbibition,  $S_{gr}$  (red + purple). Purple color highlights the intersection area of the  $H_2$ -filled pore space after drainage and after imbibition. The  $S_{gi}$  depended on the injection rate, with higher rates ( $\geq 10$  mL/h) yielding higher  $S_{gi}$ . In most cases, the  $S_{gr}$  resided in the same pores as the  $S_{gi}$  (purple) but could also redistribute to the neighboring pores indicated with red color. (For interpretation of the references to color in this figure legend, the reader is referred to the Web version of this article.)



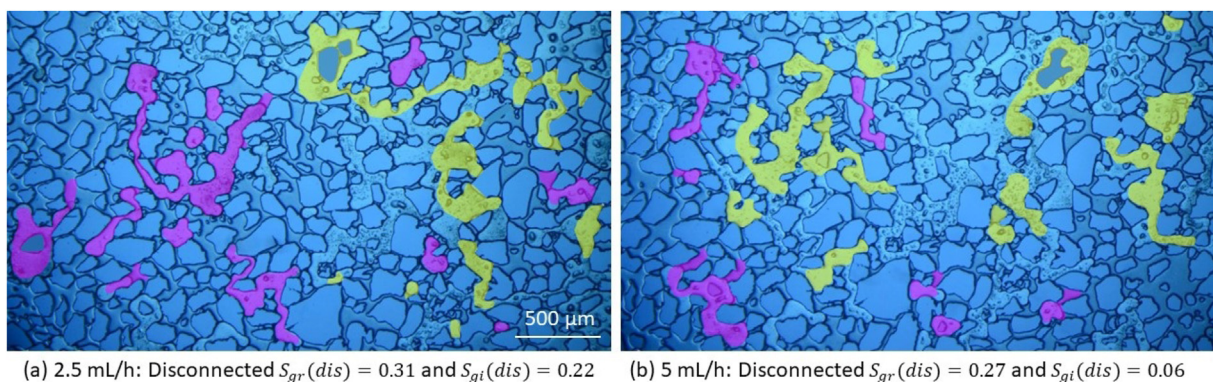
**Fig. 3** – Cyclic  $H_2$  injection and withdrawal at 5 mL/h (Exp B3), with the combined images of  $S_{gi}$  (blue + purple) and  $S_{gr}$  (red + purple). In general, the  $H_2$  distribution varied between the cycles but was similar between cycles 2 and 3. The  $S_{gi}$  tended to distribute in the large, connected pore clusters (middle right area of the images), whereas the  $S_{gr}$  distribution changed over the cycles showing hysteresis despite having similar  $S_{gr}$  values of  $\sim 0.30$ . (For interpretation of the references to color in this figure legend, the reader is referred to the Web version of this article.)

mechanisms were independent of pressure in the 5–40 bar range.

Cyclic injections resulted in fluctuating  $H_2$  saturation between 0.42 and 0.67 after drainage, due to the  $H_2$  movement

from outside the FoV (Fig. 3, Exp B3).  $H_2$  preferentially occupied the large, connected pore clusters (middle, right region of the images in Fig. 3). The residual  $H_2$  saturation after imbibition exhibited little variation over several cycles (average





**Fig. 4** – H<sub>2</sub> reconnection with the injected H<sub>2</sub> in the subsequent drainage cycle: (a) From cycle 3 to 4 at 2.5 mL/h (3 → 4), and (b) from cycle 4 to 5 at 5 mL/h (4 → 5). H<sub>2</sub> reconnection seemed somewhat stochastic locally but was favored in the large pore clusters with wide pore throats. A portion of the disconnected H<sub>2</sub> after imbibition,  $S_{gr}(dis)$  (yellow + purple) connected with the injected H<sub>2</sub> during the subsequent drainage (yellow) and the rest remained as disconnected,  $S_{gi}(dis)$  (purple). The injected H<sub>2</sub> is not manually segmented and appears in the image in its original light blue color. (For interpretation of the references to color in this figure legend, the reader is referred to the Web version of this article.)

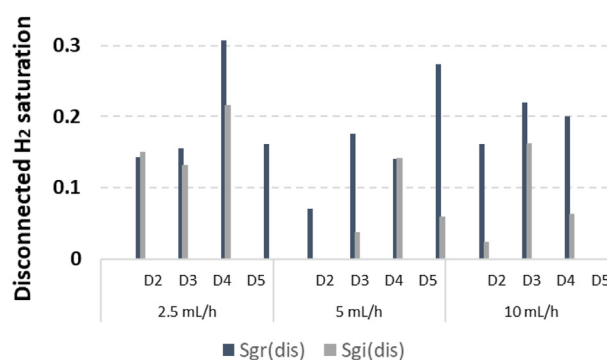
$S_{gr} = 0.31 \pm 0.03$ ), however, its pore space distribution varied between cycles due to hysteresis. Variations in the residual H<sub>2</sub> distribution have also been reported during core flooding and  $\mu$ CT imaging, despite equal residual H<sub>2</sub> saturations [30]. In our case, the residually trapped H<sub>2</sub> was not necessarily immobile in the subsequent cycle and could reconnect with the injected H<sub>2</sub>, described next.

The ability of H<sub>2</sub> ganglia to reconnect seemed stochastic locally between pores but the global distribution appeared to depend on the pore cluster morphology (Fig. 4). The H<sub>2</sub> ganglia in proximity to pore clusters with wide pore throats tended to reconnect during drainage (yellow in Fig. 4), whereas H<sub>2</sub> ganglia remained disconnected in the pores with narrow pore throats (purple in Fig. 4). In contrast, the core flooding experiments with  $\mu$ CT imaging showed that the residual CO<sub>2</sub> ganglia size gradually changed with increasing number of cycles, penetrating smaller pore throats [32,33]. Reconnection of H<sub>2</sub> ganglia during drainage was in general high, characterized by the amount reduction of H<sub>2</sub> ganglia in 9 out of 12 drainage injections relative to previous imbibition injections (Fig. 5). High H<sub>2</sub> ability to reconnect is favorable for real storage projects, reducing H<sub>2</sub> loss during cyclic injections.

### Microscopic storage capacity

Microscopic H<sub>2</sub> storage capacity was evaluated based on capillary number correlation (CNC) and pore pressures (Fig. 6, Table 3). The initial H<sub>2</sub> saturation after drainage changed both monotonically (at 5 bar) and nonmonotonically (at 1, 30 and 40 bar) with increased capillary number. A monotonic increase after a plateau region was consistent with classic CNC at core scale [44] and some microfluidic studies [41,45]. Non-monotonic trends were also reported from micromodels [46–48], likely due to the crossover from capillary to viscous flow regimes and/or micromodel properties [44].

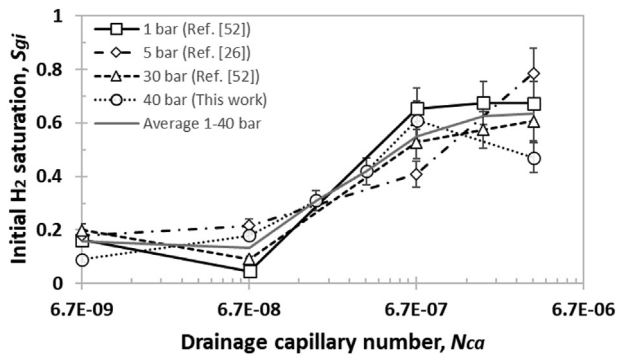
When averaged for a specific capillary number, the initial H<sub>2</sub> saturation exhibited a monotonic trend (Fig. 6). The critical drainage capillary number ranged between  $3.4\text{--}6.8 \times 10^{-7}$ , corresponding to maximum H<sub>2</sub> storage capacity of ~60% of the



**Fig. 5** – Quantification of the H<sub>2</sub> ganglia reconnection, by comparing the disconnected H<sub>2</sub> saturation after imbibition,  $S_{gr}(dis)$  (blue) and the subsequent drainage cycle,  $S_{gi}(dis)$  (gray) at the injection rates of 2.5, 5 and 10 mL/h. The horizontal axis compares the two subsequent cycles: The transition from the imbibition cycle 1 to the drainage cycle 2 is denoted as D2. In most cases, the disconnected H<sub>2</sub> saturation decreased in the subsequent drainage cycle, indicating high reconnection ability. (For interpretation of the references to color in this figure legend, the reader is referred to the Web version of this article.)

pore space. The optimal reservoir scale injection rate yielding the highest storage capacity would therefore be in the range of ~170–340 thousand standard m<sup>3</sup>/day, assuming the injector perforation length of 30 m and the experimental injection velocity. Our storage capacity and injection rates were comparable with the reservoir simulations of aquifer storage assuming maximum H<sub>2</sub> saturations of 70% and injection rates of ~200–300 thousand standard m<sup>3</sup>/day [11,12,49].

No clear pressure effect was observed on the initial H<sub>2</sub> saturation, contradicting classic threshold pressure phenomena [50]. The saturation independence from pressure was likely due to insignificant wettability and interfacial tension alterations in the H<sub>2</sub>–H<sub>2</sub>O systems under the studied pressure range of 1–40 bar. The H<sub>2</sub> contact angles (i.e. wettability)



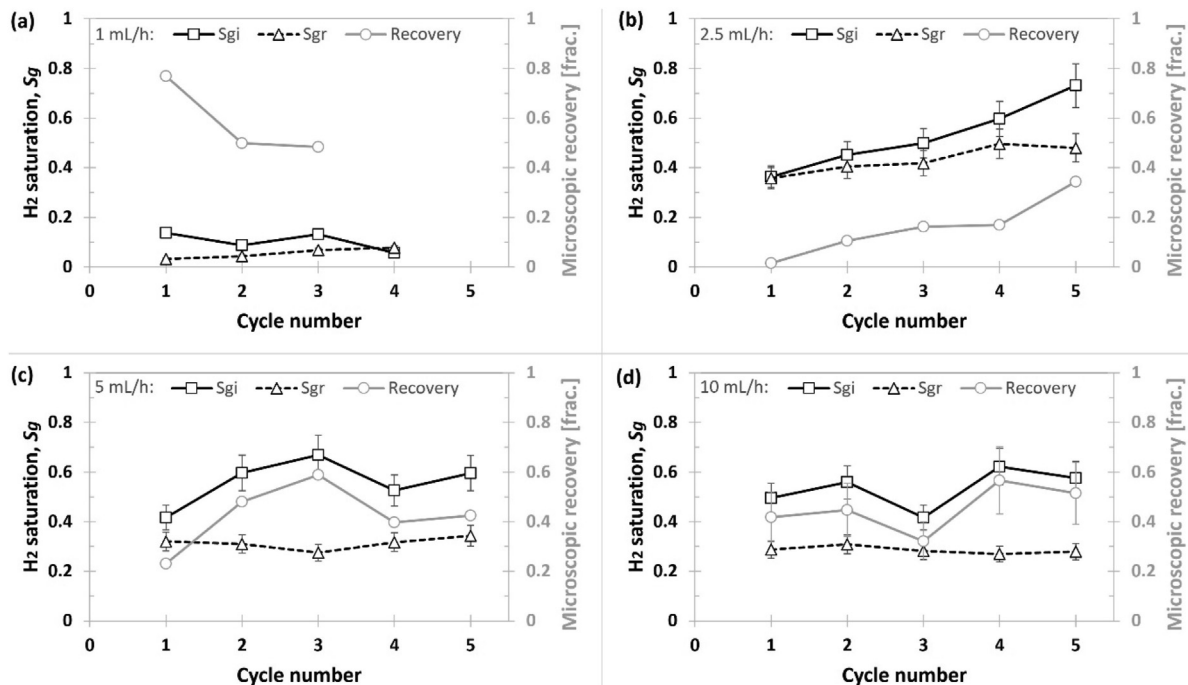
**Fig. 6** – Microscopic H<sub>2</sub> storage capacity from the capillary number correlation (CNC) compared with the results from the same micromodel at 1, 5, and 30 bar [26,52]. The  $S_{gi}$  was independent of pressure. The maximum H<sub>2</sub> storage capacity was on average equal to ~60% of the pore space (gray curve). The error bars represent the image analysis relative uncertainty (9%).

showed no pressure dependence under the range of 20–100 bar in Berea and Bentheimer sandstones [14]. Insignificant contact angle changes of  $\sim 5^\circ$  were reported for H<sub>2</sub> on the pure quartz surface for pore pressures ranging between 1 and 50 bar and room temperature [16]. The H<sub>2</sub>–H<sub>2</sub>O interfacial tension increases by less than 1% from 1 to 40 bar [38]. No correlation between initial gas saturation and pressure has been reported for CO<sub>2</sub>, N<sub>2</sub>, and H<sub>2</sub> at core scale [30,51], corroborating our results.

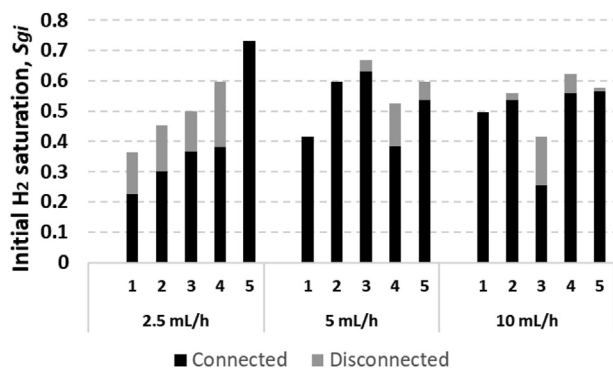
### Residual trapping and microscopic recovery during cyclic injections

The initial and residual H<sub>2</sub> saturations were quantified for cyclic injections and the corresponding microscopic recovery factor was calculated for every cycle (Fig. 7, Table 3). The H<sub>2</sub> saturation range between the cycles depended on the injection rate, in accordance with the critical capillary number from the CNC (Fig. 6). The low injection rate (1 mL/h) resulted in a low saturation range between 0.03 and 0.14 (Fig. 7a), with an increase up to 0.27–0.73 at higher injection rates ( $\geq 2.5$  mL/h) (Fig. 7b–d). The initial H<sub>2</sub> saturation varied between the cycles at higher injection rates, whereas the residual H<sub>2</sub> remained nearly constant and equal to  $\sim 0.43$  at 2.5 mL/h,  $\sim 0.31$  at 5 mL/h, and  $\sim 0.29$  at 10 mL/h, explained next.

The reason for the fluctuations in the initial H<sub>2</sub> saturation could be twofold: 1) H<sub>2</sub> redistribution from outside the FoV, caused by random H<sub>2</sub> injection patterns, and/or 2) the presence of the disconnected H<sub>2</sub> ganglia. A further analysis (Fig. 8) revealed that both connected and disconnected H<sub>2</sub> saturation were stochastic without any clear trend, meaning that H<sub>2</sub> redistribution (reason 1) was the main cause for the fluctuating initial saturations. The opposite was observed for *n*-hexane (C<sub>6</sub>H<sub>14</sub>, used as a proxy for CO<sub>2</sub>) cyclic injections in a micromodel, where the amount of the disconnected C<sub>6</sub>H<sub>14</sub> increased over the cycles due to the converged injection patterns through the most accessed pore channels [53]. The discrepancies with our study were likely caused by the differences in the micromodel design. In our case, two open



**Fig. 7** – Initial ( $S_{gi}$ ) and residual ( $S_{gr}$ ) H<sub>2</sub> saturation and the resulting microscopic recovery factors during cyclic injections at the flow rate of: (a) 1 mL/h, (b) 2.5 mL/h, (c) 5 mL/h, (d) 10 mL/h. The  $S_{gi}$  fluctuated between the cycles due to H<sub>2</sub> redistribution from outside the FoV. In contrast, the  $S_{gr}$  showed better reproducibility, with nearly constant values throughout the cycles. The microscopic recovery factors reflected the fluctuations in the  $S_{gi}$ . The error bars represent the image analysis relative uncertainty (9%).

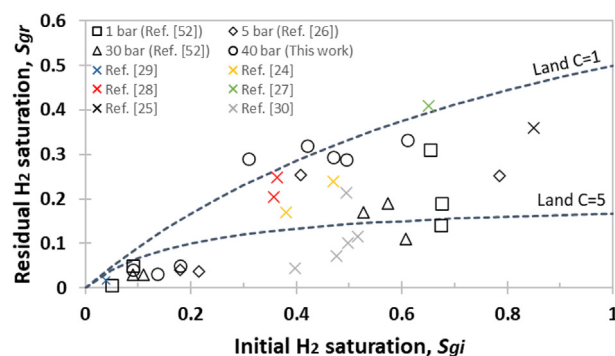


**Fig. 8** – Share of the connected and disconnected initial  $H_2$  saturation ( $S_{gi}$ ) during cyclic injections (cycles 1–5) at the injection rate of 2.5, 5 and 10 mL/h. No clear trend was observed, meaning that the presence of the disconnected  $H_2$  did not cause the fluctuations in the  $S_{gi}$  in Fig. 7.

channels along the micromodel length (Fig. 1) resulted in crossflow, facilitating more random injection patterns in multiple directions. In the case of  $C_6H_{14}$  injections, the open channels were built in the opposite direction, that is along the micromodel widths, creating one-directional injection pattern.

Contrary to the initial  $H_2$  saturation, the residual  $H_2$  saturation was more reproducible because the  $H_2O$  injection was eased in a strongly hydrophilic system, with well-established injection patterns through the wetting  $H_2O$  films coating the grain surfaces. Note that natural reservoirs contain organic-rich material, making their rock surfaces more hydrophobic than our micromodel [13,16]. Greater reproducibility of the residual  $H_2$  saturation throughout the cycles is beneficial from the storage perspective, permitting a better control of the  $H_2$  loss due to residual trapping. The distribution of the residual  $H_2$ , however, visually changed over the cycles as mentioned in section 3.1 (Fig. 3). Such hysteretic behavior due to residual trapping can affect the imbibition relative permeability, and hysteresis in  $H_2$ – $H_2O$  relative permeability has already been demonstrated at core scale [24,25].

The microscopic recovery factors, defined as  $(S_{gi} - S_{gr})/S_{gi}$ , fluctuated between the cycles, in alignment with the initial  $H_2$  saturation (Fig. 7). The recovery factors ranged between 2% and 77% with an average of ~40%, comparable with the recovery factors from reservoir simulations of  $H_2$  storage in the  $H_2O$  zone of a depleted hydrocarbon field: 49% [54] and aquifer storage: 36–59% [11,55]. Note that recovery factors from reservoir simulations are macroscopic and valid for the entire reservoir, contrary to microfluidics which deal with the microscopic phenomena. The highest  $H_2$  saturation is expected in the near-well area, with gradually decreasing  $H_2$  saturation when approaching the  $H_2$ – $H_2O$  transition zone in the far-well area [11,12,49]. During cyclic injections, the  $H_2$ – $H_2O$  transition boundary is constantly moving, resulting in increasing  $H_2O$  saturation during  $H_2$  withdrawal with associated residual  $H_2$  trapping. Hence, the reported microscopic recovery factors are mostly relevant for the  $H_2$ – $H_2O$  transition zone.



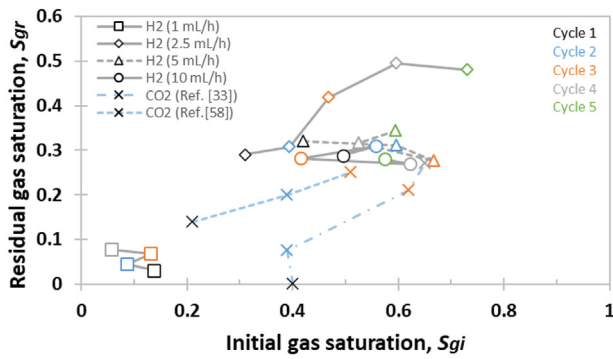
**Fig. 9** – Trapping model based on  $H_2$  saturations after primary drainage ( $S_{gi}$ ) and imbibition ( $S_{gr}$ ), combined with the results from the same micromodel at 1, 5 and 30 bar [26,52] and the literature  $H_2$  data at core scale denoted by colored crosses [24,25,27–30]. Most of the measurements followed the Land trapping model with the trapping coefficients  $C = 1$  and  $C = 5$ . The upper limit data points were comparable to  $CO_2$  with the trapping coefficients between 0.2 and 2.1 in sandstones [56,57].

#### Trapping model

We combine  $H_2$  saturations after primary drainage and imbibition together with the results from the same micromodel at 1–30 bar and available  $H_2$  data at core scale (Fig. 9), to construct  $H_2$  trapping relationship based on an empirical Land model [34]. This model was derived from the measurements of the initial and residual gas saturations in sandstone core samples, defined as follows:  $S_{gr} = S_{gi}/(1 + C \cdot S_{gi})$ , where  $C$  is the trapping coefficient. The data points were greatly scattered and mostly fell within the trapping coefficient ( $C$ ) range between 1 and 5. The upper boundary points matched the  $CO_2$  trapping models in different sandstones, with the trapping coefficient range of 0.2–2.1 [56,57]. In contrast, the lower boundary points were outside the reported  $CO_2$  data, indicating that less  $H_2$  trapping may be expected compared to  $CO_2$ . The lower boundary points disagreed with the contact angle measurements which reported less  $H_2$  wetting (more hydrophilic) compared to  $CO_2$  [13,16], that in turn implies more  $H_2$  trapping [23]. Greater scatter and disagreement of lower boundary  $H_2$  data points with  $CO_2$  data could be due to the differences in the porous materials and methodologies. The  $CO_2$  measurements were obtained from conventional core flooding with the core length of 6–12 cm, whereas microfluidics and shorter core plugs of 1.5–5.7 cm were used for most  $H_2$  saturation measurements.

The initial and residual  $H_2$  saturations from cyclic injections (Fig. 10) showed no significant increase in residual trapping over the cycles for similar initial  $H_2$  saturations, consistent with the Land model and corroborated by  $H_2$  [30] and some  $CO_2$  cyclic injection studies [58,59]. In contrast, some  $CO_2$  measurements deviated from the Land model, with a sharp increase in residual saturations over the injection cycles [31–33]. The exact mechanism for this deviation is still poorly understood but could be attributed to pore throat blockage due to fines migration,  $CO_2$  adhesion to the grain surfaces, and/or wettability alteration to a “patchy” mixed-





**Fig. 10** – Initial ( $S_{gi}$ ) and residual ( $S_{gr}$ )  $H_2$  saturations during cyclic injections at the injection rates of 1, 2.5, 5 and 10 mL/h. The  $S_{gr}$  were nearly stable over the injection cycles for similar  $S_{gi}$ , consistent with the Land trapping model. Our results were compared with  $CO_2$  studies (denoted by crosses), where there is currently a contradiction, with some studies following the Land model [58] and other studies deviating from it [33].

wet with discontinuous  $CO_2$ -wet areas [33,60]. Even though our results and one core scale  $H_2$  study [30] did not indicate a significant increase in residual trapping, lack of  $H_2$  studies and the disagreement in the  $CO_2$  literature emphasizes the importance for further investigations of  $H_2$  cyclic injections. A potential increase in residual  $H_2$  trapping over the injection cycles is undesired as it will reduce  $H_2$  storage efficiency.

Trapping models based on the 2D microfluidic experiments have a limited applicability for 3D reservoirs due to the small volume and the absence of gravity and heterogeneity. The same applies for the storage capacities and recovery factors quantified in sections 3.2 and 3.3. The 2D micromodels are suitable for qualitative description of the pore scale flow mechanisms, which can support core scale measurements. Extrapolation of the quantitative results to natural reservoirs requires caution and should be preferably done by pore scale modelling. However, our results followed classic CNC trends (Fig. 6) and the  $H_2$  saturations were within the literature range at core scale (Fig. 9). Therefore, it is reasonable to claim that the quantified  $H_2$  storage capacities and trapping model from microfluidics can temporarily substitute for missing measurements until a comprehensive core scale dataset is available.

## Conclusions

We report a series of cyclic  $H_2$ – $H_2O$  injections in a micro-model mimicking sandstone pore patterns and relevant for shallow  $H_2$  storage in aquifers and depleted gas fields with an underlying water zone. We found that  $H_2$  saturation after primary drainage increased with increasing capillary number, with maximum storage capacities up to ~60% of the pore space. When combined with previous results from the same micromodel, the initial  $H_2$  saturation was independent of pressure in the range of 1–40 bar. The distribution of initial and residual  $H_2$  in the pore space were hysteretic over the injection cycles, with fluctuating initial but similar residual  $H_2$

saturations. The residually trapped  $H_2$  showed good reconnection ability, which was favored in proximity to the large pore clusters with wide pore throats. The  $H_2$  trapping followed the Land model, with trapping coefficient between 1 and 5 where the upper limit values matched the  $CO_2$  trapping models in sandstones. The microscopic  $H_2$  recovery factors varied due to the fluctuating initial  $H_2$  saturation and on average were equal to ~40%, relevant for the  $H_2$ – $H_2O$  transition zone in the far-well area. Higher reconnection ability and reproducibility of residual  $H_2$  saturation are beneficial for underground  $H_2$  storage but this positive impact may be suppressed by its hysteretic distribution over the injection cycles. Future work should be focused on core scale cyclic injections and on pore scale modelling for upscaling to natural reservoirs.

## Declaration of Competing Interest

The authors declare that they have no known competing financial interests or personal relationships that could have appeared to influence the work reported in this paper.

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